

Topological Gravity and Matrix Models

Robbert Dijkgraaf

Institute for Advanced Study

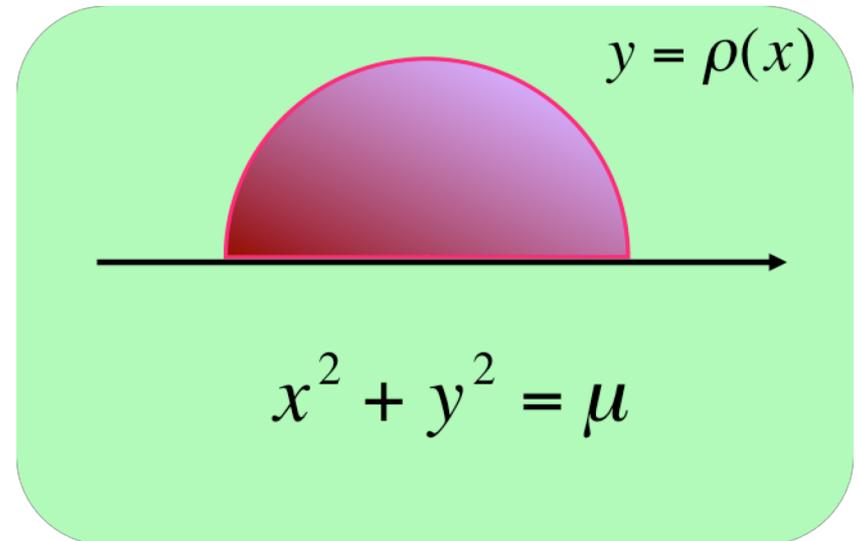
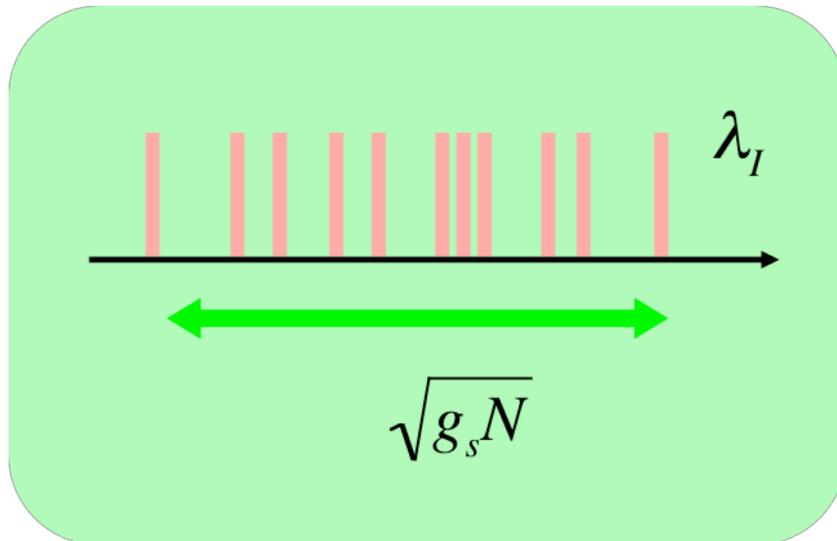
Based on work with E. Witten, arXiv:1804.03275 and in progress

Wigner's Random Matrix Model

$$\lim_{N \rightarrow \infty} Z_N, \quad Z_N = \frac{1}{\text{Vol } U(N)} \int_{N \times N} d\Phi \cdot e^{-\text{Tr } \Phi^2 / g_s}$$

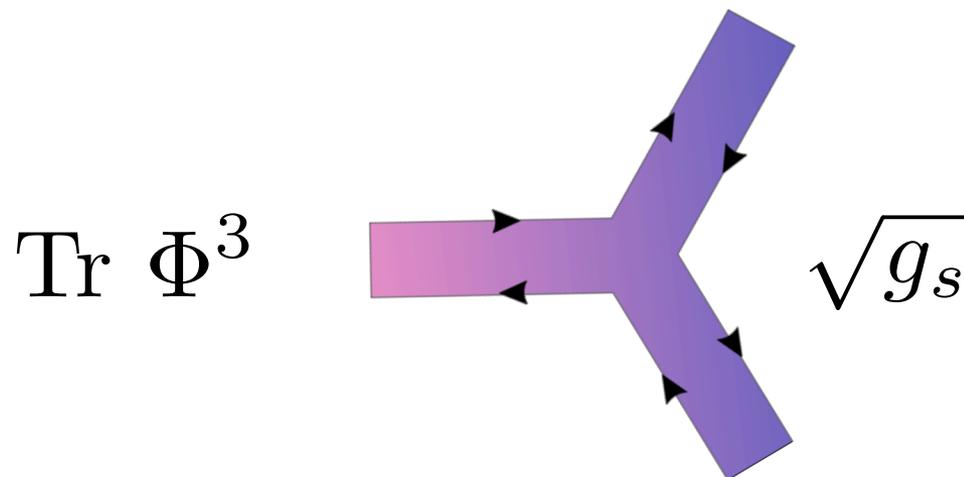
Eigenvalue distribution in 't Hooft limit

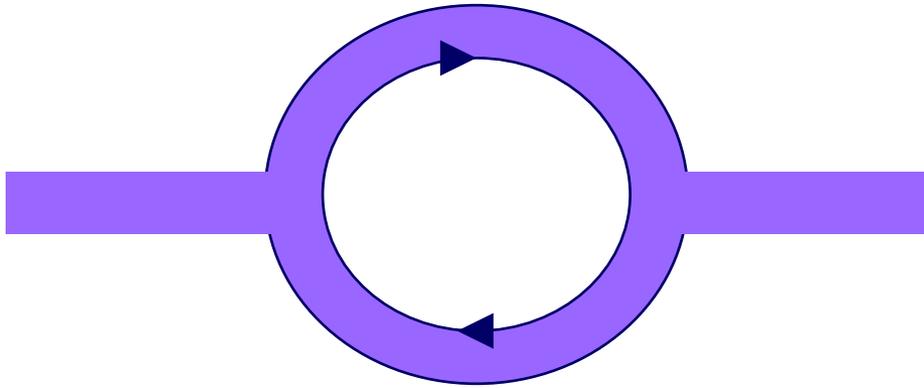
$$N \rightarrow \infty, \quad g_s \rightarrow 0, \quad Ng_s = \mu = \text{fixed}$$



Matrix Models and 2d Gravity

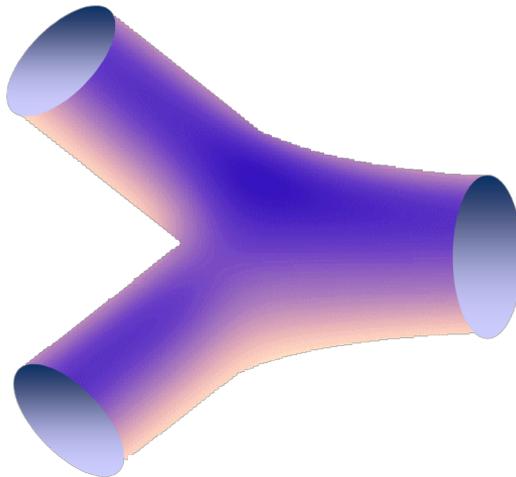
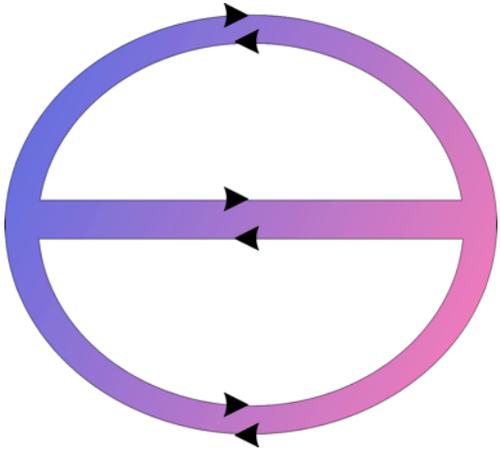
$$\int d\Phi e^{-\frac{1}{g_s} \text{Tr}(\Phi^2 + \Phi^3)}$$





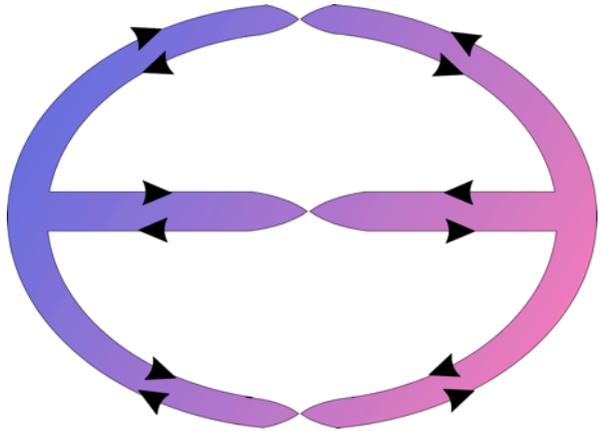
$$g_s N = \mu = \textit{fixed}$$

Planar diagrams



$$N^2$$

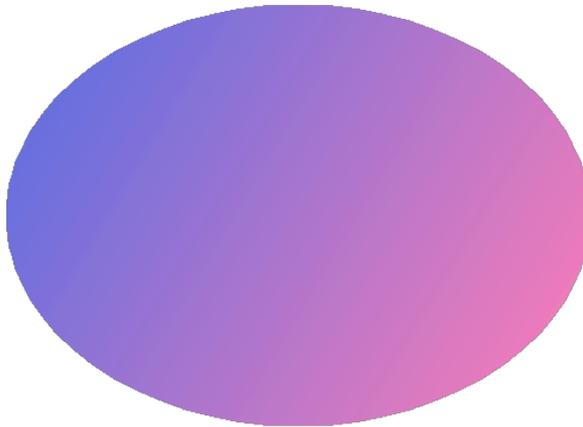
Non-planar diagrams



$$g = 1$$

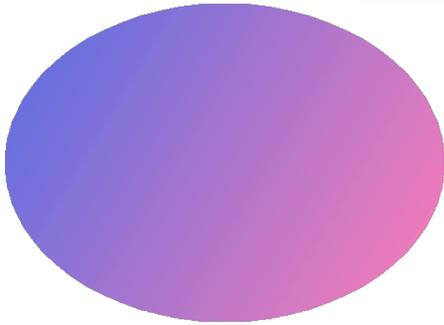
$$N^{2-2g}$$

Sum of diagrams

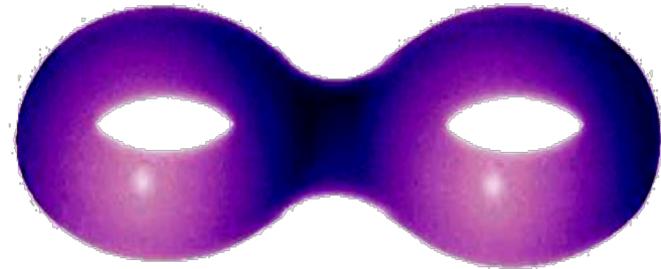


String Partition Function

$$Z = \exp \sum_{g \geq 0} g_s^{2g-2} F_g(\mu)$$

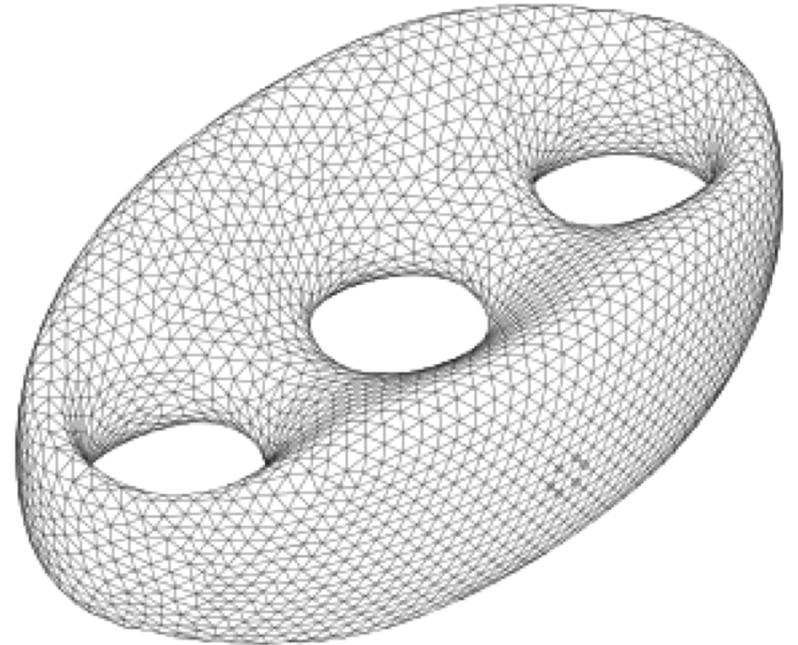
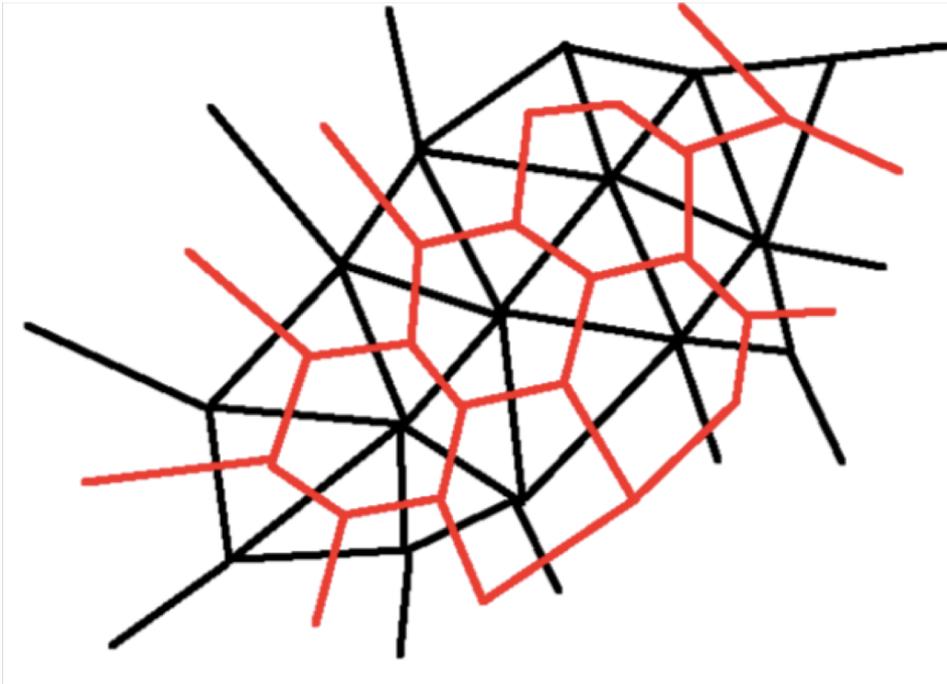


ribbon diagrams
1/N expansion



genus g
string surface

Triangulations of surface



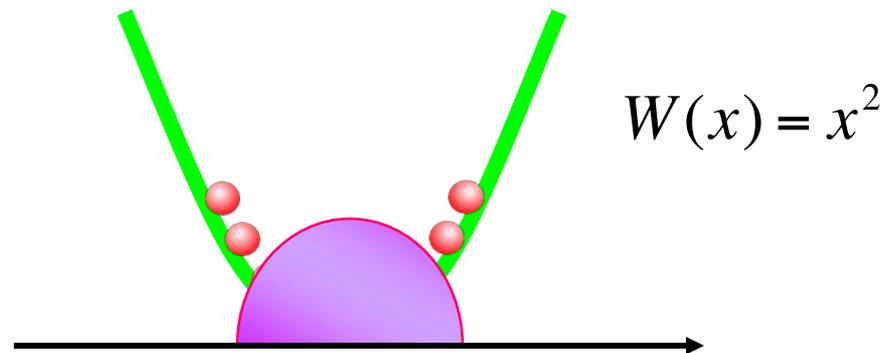
Dual graph
Limit of large number of vertices

Eigenvalue Dynamics

$$Z_{matrix} = \int d^N \lambda \cdot \prod (\lambda_I - \lambda_J)^2 \cdot e^{-\sum_I \lambda_I^2 / g_s}$$

Effective action (repulsive Coulomb force)

$$S_{eff} = \sum_I \lambda_I^2 - 2g_s \sum_{I < J} \log(\lambda_I - \lambda_J)$$



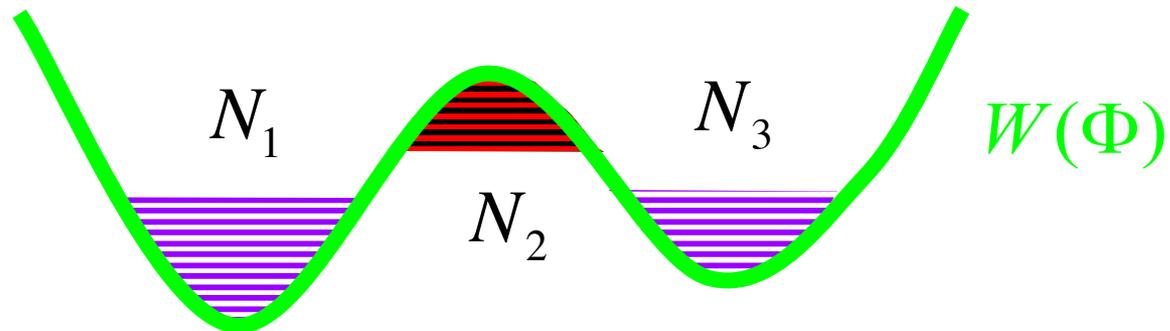
General Matrix Model

$$Z_{matrix} = \frac{1}{\text{Vol } U(N)} \int_{N \times N} d\Phi \cdot e^{\text{Tr } W(\Phi)/g_s}$$

't Hooft limit

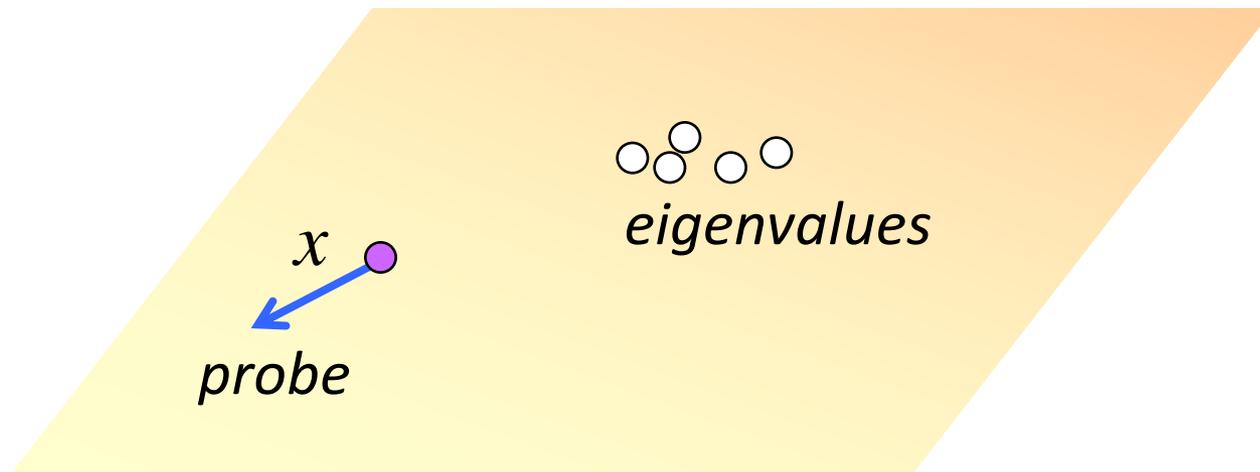
$$N_I \rightarrow \infty, \quad g_s \rightarrow 0, \quad N_I g_s = \mu_I = \text{fixed}$$

Filling fractions (perturbative expansion)



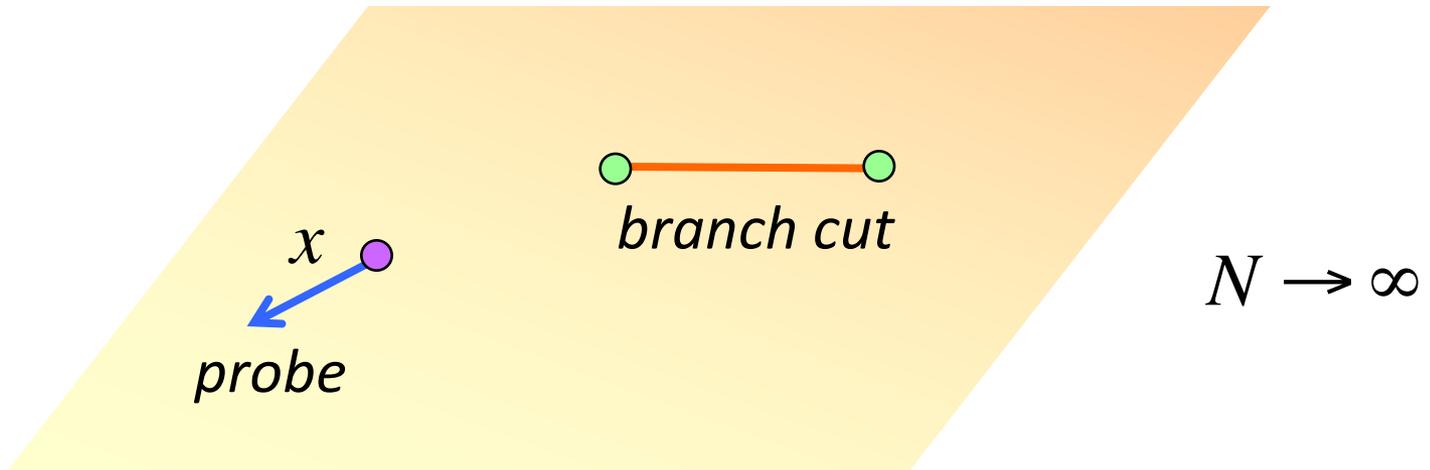
Resolvent

$$y dx = dS_{\text{eff}} = dW(x) + 2g_s \text{Tr} \frac{dx}{\Phi - x}$$



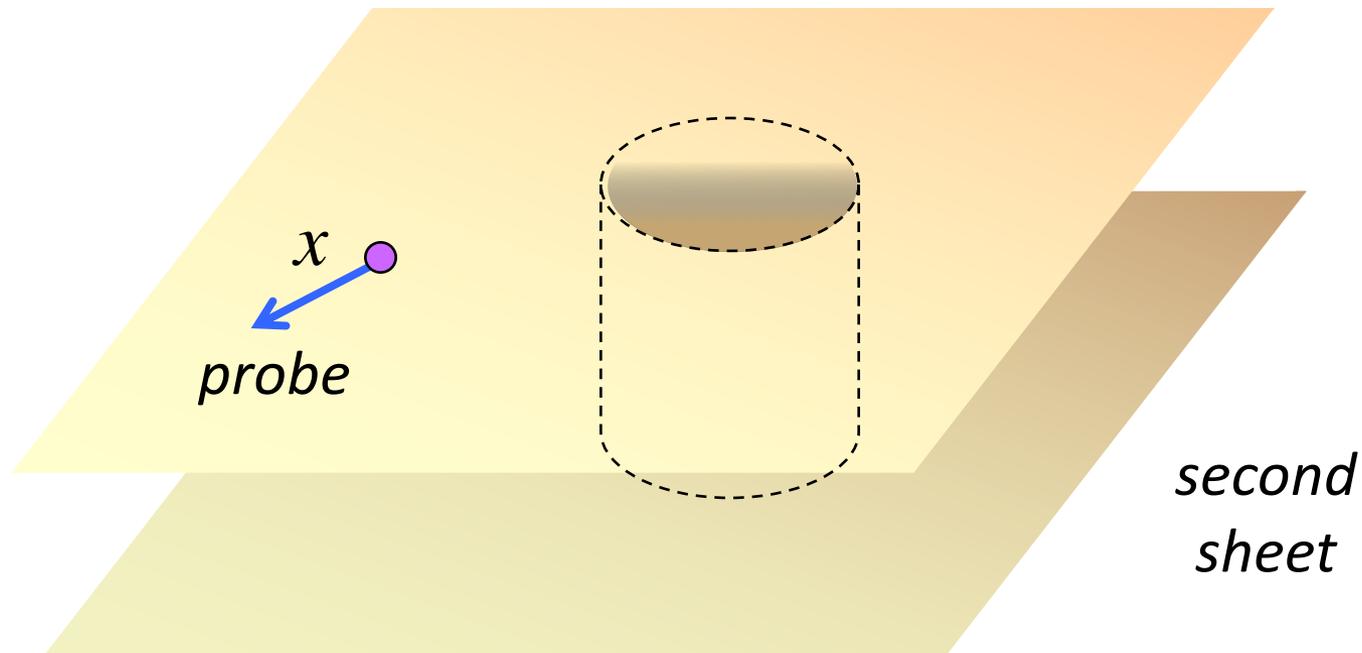
Resolvent

$$y dx = dS_{\text{eff}} = dW(x) + 2g_s \text{Tr} \frac{dx}{\Phi - x}$$



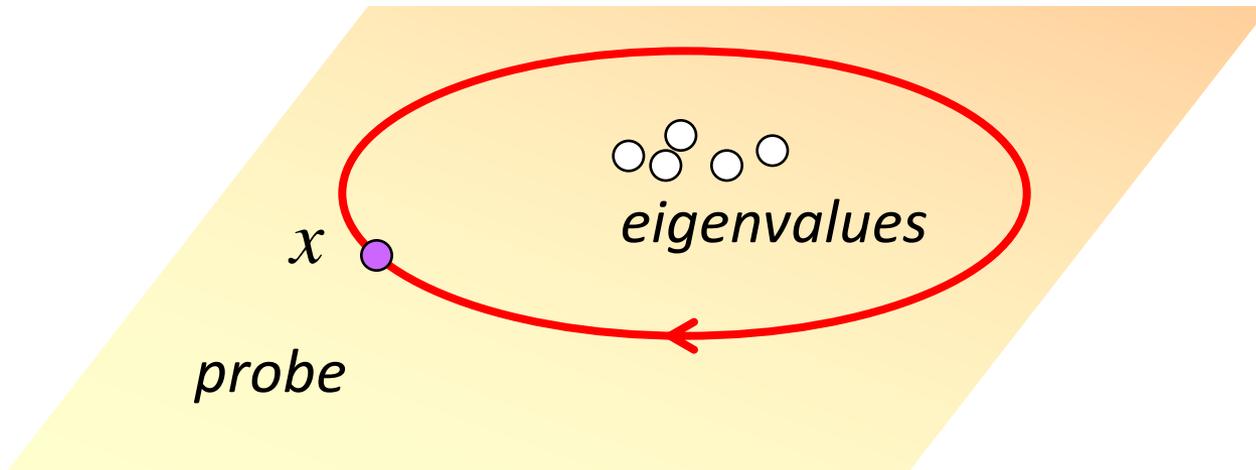
Effective Geometry

$$y^2 = W_n'(x)^2 + f_{n-1}(x) = P_{2n}(x)$$



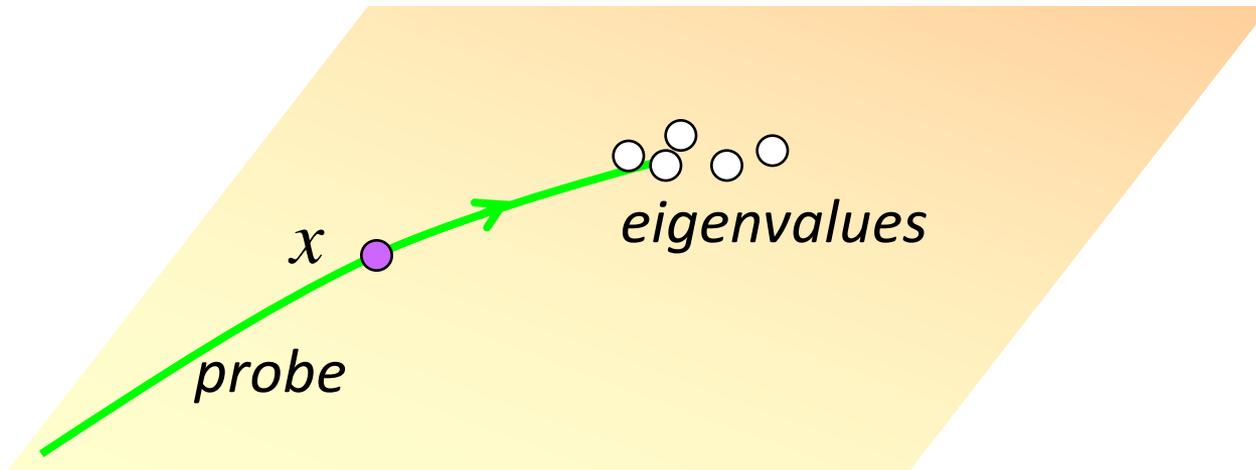
Periods (A)

$$\oint_{A_I} y dx = g_s N_I = \mu_I$$



Periods (B)

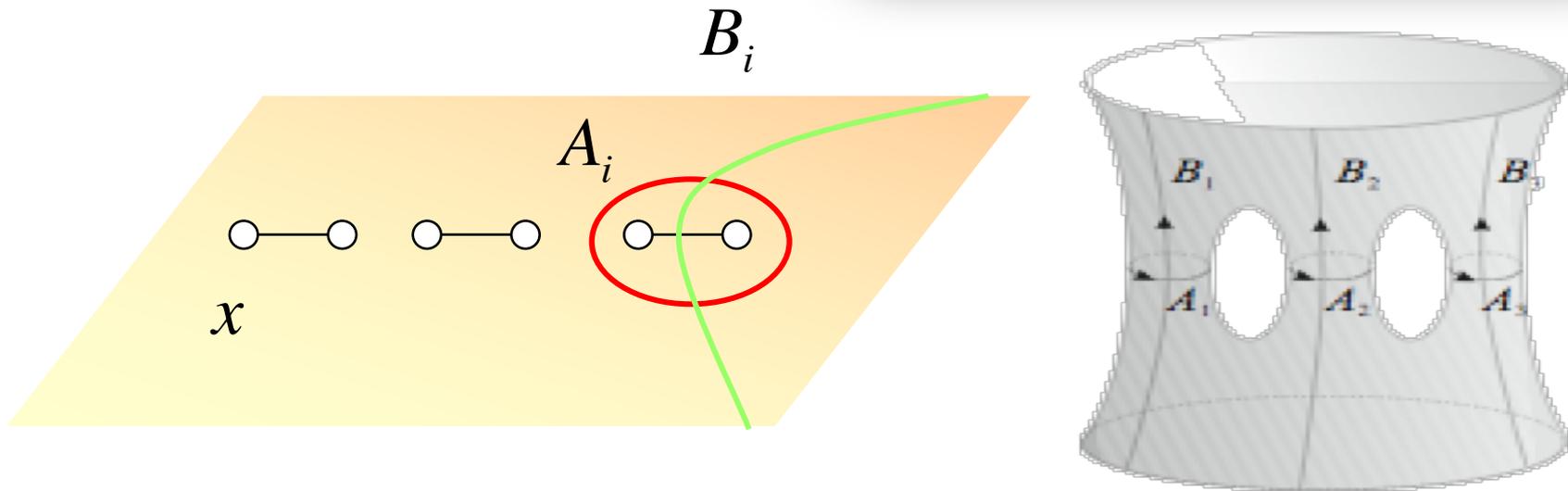
$$\int_{B_I} y dx = \frac{\partial F}{\partial \mu_I}$$



Spectral Curve

Hyperelliptic curve

$$\Sigma: y^2 = W'(x)^2 + f(x)$$



Quantum invariants of Σ , $\omega = ydx$

$$\mu_i = \oint_{A_i} ydx$$

$$\frac{\partial F_0}{\partial \mu_i} = \oint_{B_i} ydx$$

Spectral Curve

General phenomenon in matrix models (chains, external fields, beta-ensembles, SO/Sp, etc). Spectral curve

$$P(x, y) = 0$$

With meromorphic one-form

$$\omega = ydx$$

Emergent geometry of large N systems

Complex Curves in Phase Space

algebraic curve = level set

$$\Sigma : H(x, y) = 0$$

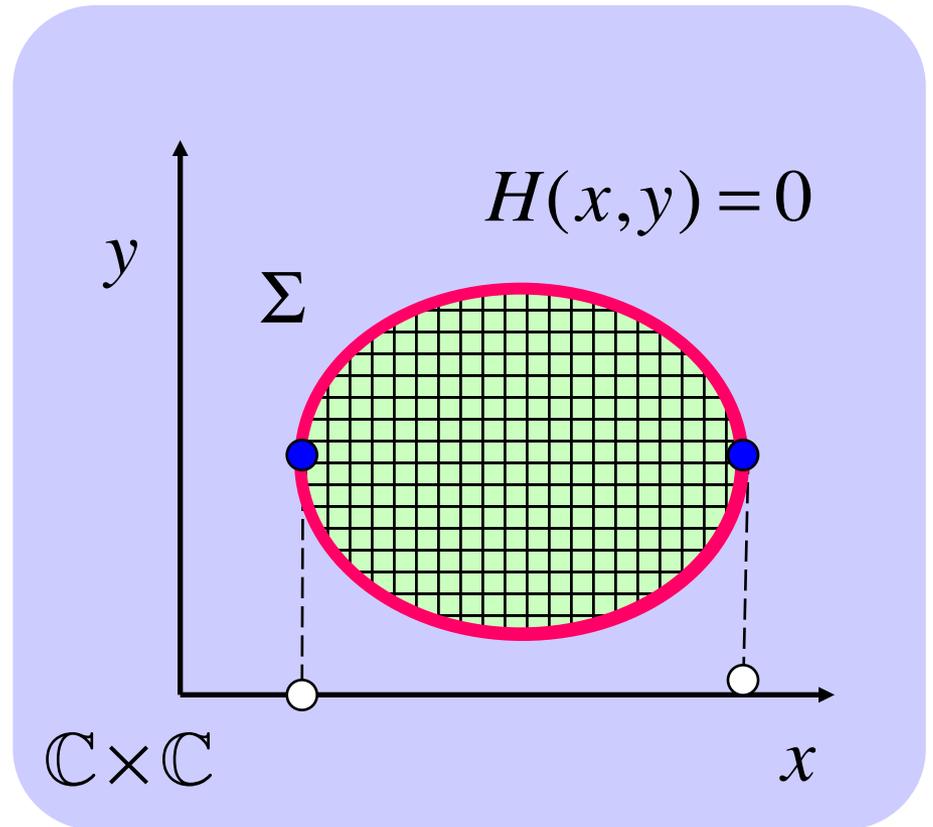
Hamilton-Jacobi theory

$$y = p(x)$$

branch pts = turning pts

Liouville form $\omega = ydx$

action
$$S(x) = \int_{x_*}^x \omega$$



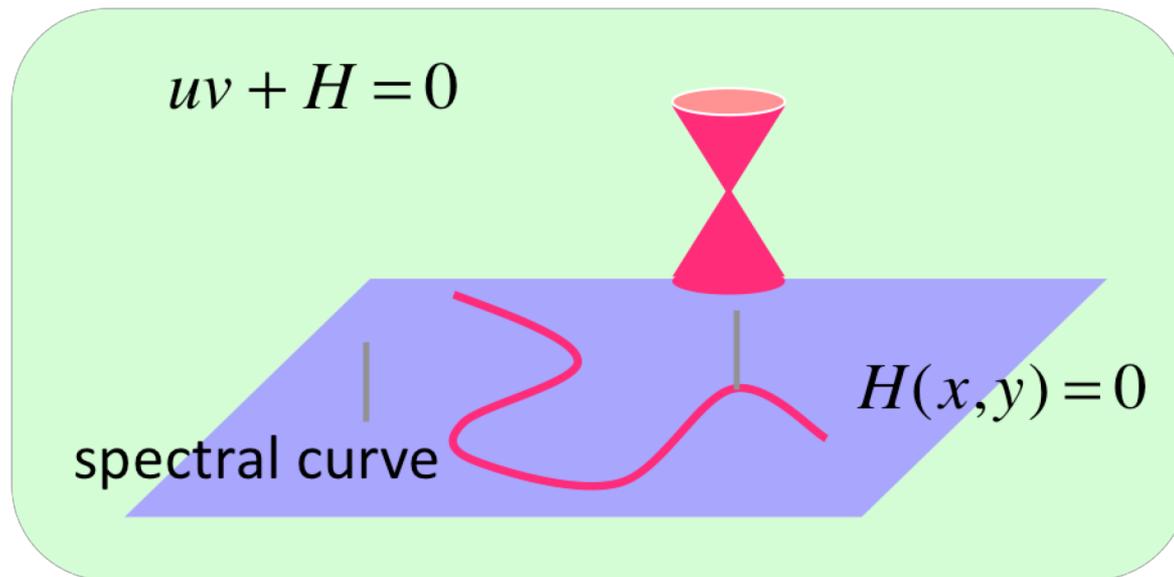
BS Quantization

$$dx \wedge dy \sim \hbar = g_s$$

Topological string on hypersurface CY_3

Calabi-Yau hypersurface in \mathbb{C}^4

$$X : uv + H(x, y) = 0$$

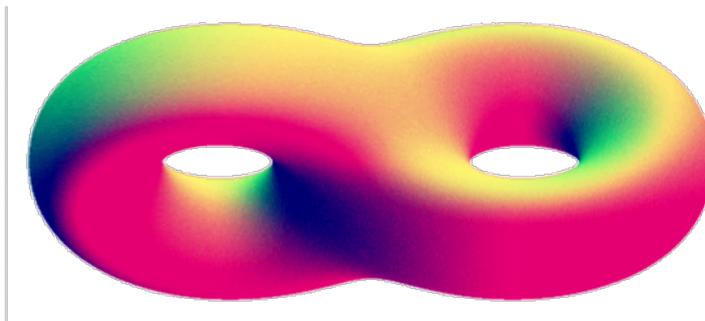


Period maps

$$\Omega = \frac{du}{u} \wedge dx \wedge dy \rightarrow \omega = ydx$$

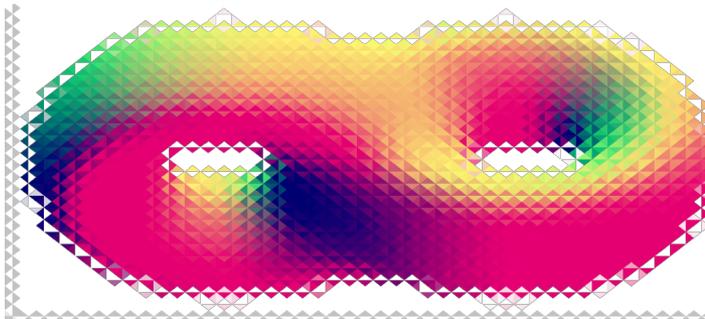
Quantum Curves

$N \rightarrow \infty$



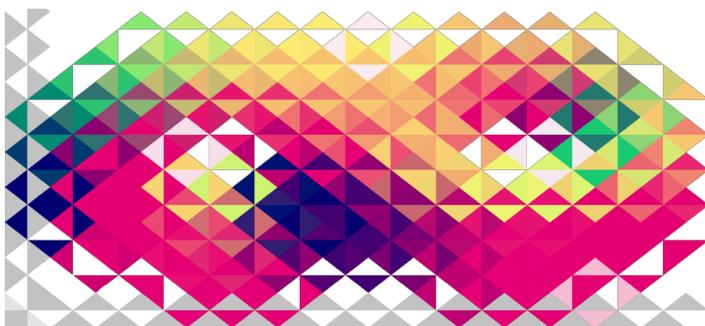
$$F_0(\mu)$$

$O(1/N)$



$$\sum_{g \geq 0} N^{2-2g} F_g(\mu)$$

N finite



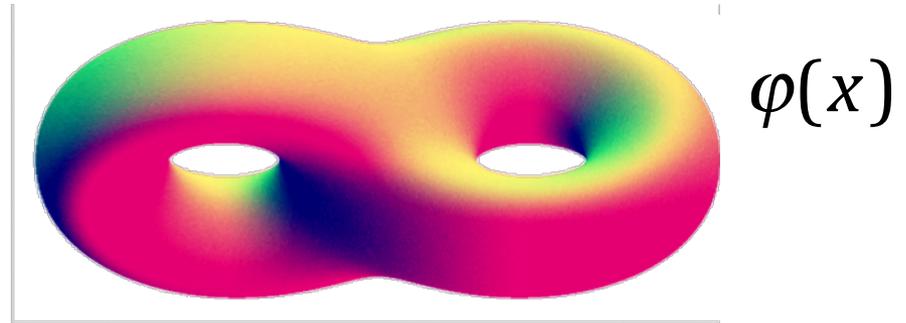
non-perturbative

Matrix models and chiral CFT

Eigenvalue density/matrix resolvent = collective boson field

$$\partial\varphi(x) = \text{Tr} \frac{1}{x - \Phi}$$

Free scalar field



$$Z_{matrix} = \int D\varphi \cdot e^{-S[\varphi]}, \quad S = \int (\partial\varphi)^2 + O(g_s)$$

Observables

$$\partial\varphi = \sum \text{Tr} \Phi^n x^{-n-1} \quad S + \sum t_n \text{Tr} \Phi^n$$

Loop Equations & Virasoro Constraints

Invariance under diffeomorphism $x \rightarrow \tilde{x}(x)$

Generated by stress-tensor

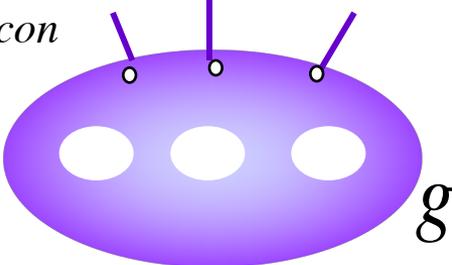
$$T(x) = \frac{1}{2}(\partial\phi)^2(x)$$

Loop equations

$$\langle T(x) \rangle = W'(x)^2 + f(x)$$

Eynard-Orantini: All-Genus Solution

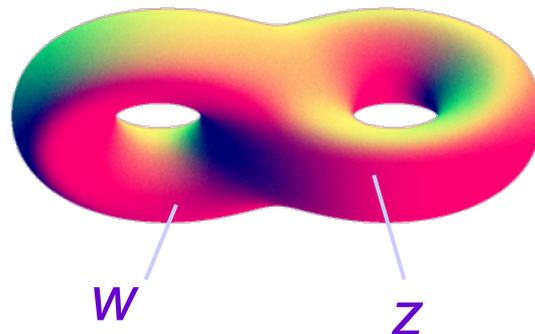
Recursion relations for correlation functions, completely geometric

$$\begin{aligned}
 W(x_1, \dots, x_n) &= \left\langle \text{Tr} \frac{1}{x_1 - \Phi} \dots \text{Tr} \frac{1}{x_n - \Phi} \right\rangle_{con} \\
 &= \langle \partial\varphi(x_1) \dots \partial\varphi(x_n) \rangle_{con} \\
 &= \sum_g g_s^{2g-2} W_g
 \end{aligned}$$


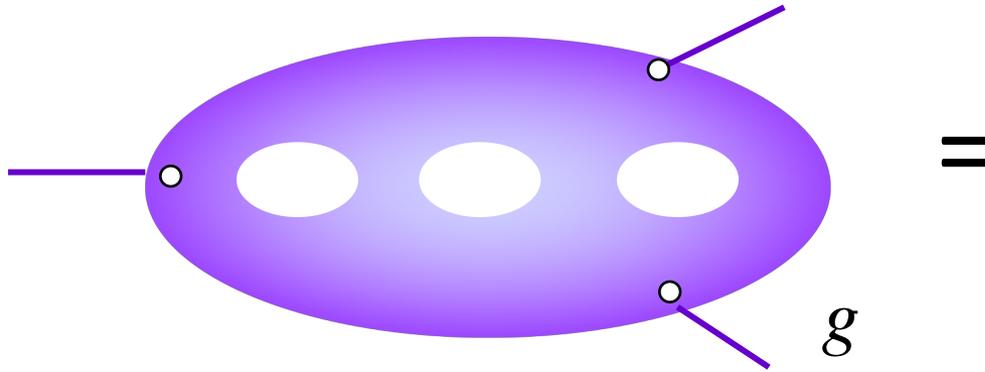
A purple genus-3 surface (a torus with three holes) is shown. Three small white circles are marked on the top boundary, each with a purple line pointing upwards, representing insertion points for the correlation function.

Lowers genus, adds more insertions. Reduces to Bergmann kernel

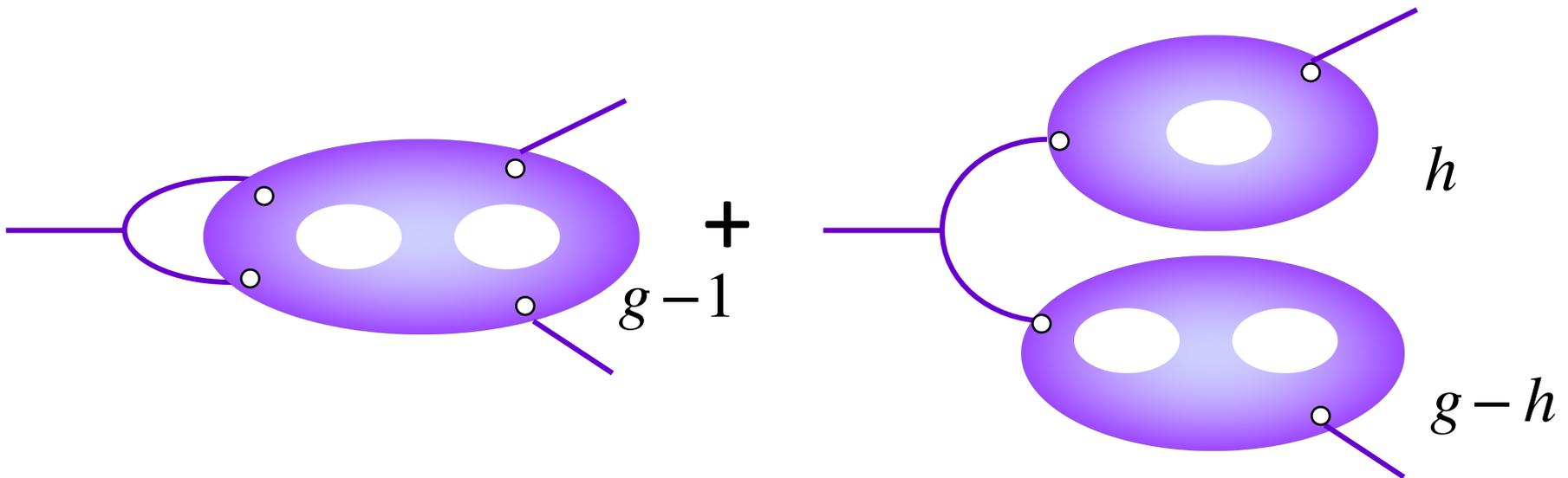
$$W_0(w, z) = B(w, z) dw dz$$



Recursion Relations

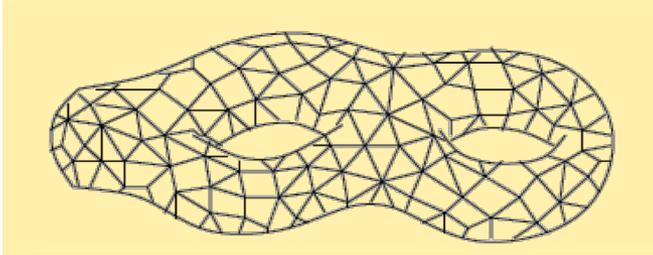


=



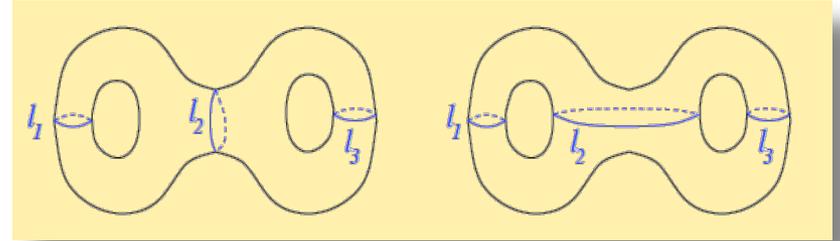
Quantum Spectral Curves

Random surfaces



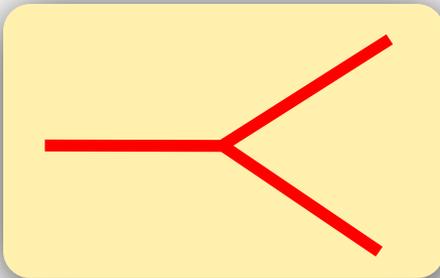
$$y^p = x^q$$

Geometry of \mathcal{M}_g



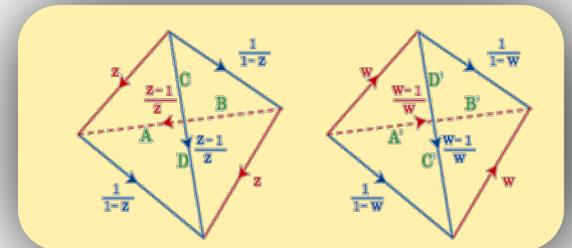
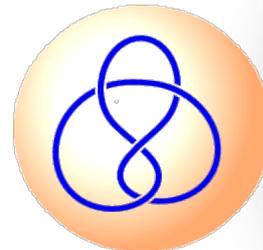
$$y = \sin \sqrt{x}$$

Toric Calabi-Yau



$$e^x + e^{-y} = 1$$

Hyperbolic knots



$$A(e^x, e^y) = 0$$

Double Scaling Limits

Random triangulations: 2d gravity coupled to (p,q) minimal CFT

$$P(x, y) = 0 \longrightarrow y^p = x^q + \dots$$

Pure 2d gravity $(p,q) = (2,3)$

$$S = \int \sqrt{g}R + \mu\sqrt{g} \qquad y^2 = x^3$$

Topological gravity $(p,q) = (2,1)$

$$y^2 = x$$

Topological Gravity

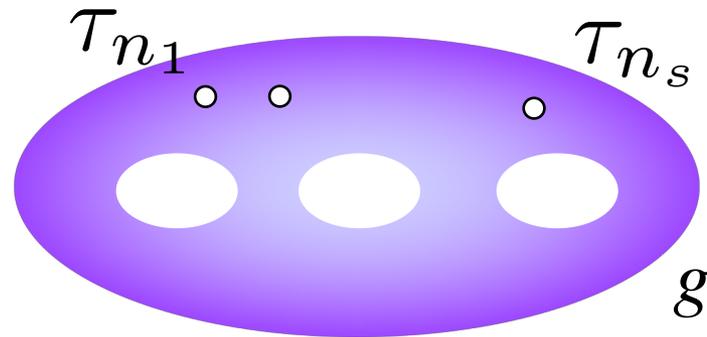
Moduli space of Riemann surfaces

$$\mathcal{M}_{g,s}, \quad \dim \mathcal{M}_{g,s} = 3g - 3 + s$$

Observables

$$\tau_n = c_1(L)^n$$

Correlation functions



$$\langle \tau_{n_1} \cdots \tau_{n_s} \rangle_g = \int_{\mathcal{M}_{g,s}} c_1(L_1)^{n_1} \wedge \cdots \wedge c_1(L_s)^{n_s}$$

Maryam Mirzakhani's Work

Compute the volume of moduli space of surfaces

$$\mathcal{F}_g = \text{vol}(\mathcal{M}_g)$$

Weil-Peterson metric: pick a constant negative curvature metric on the surface, equivalent to flat $SL(2, \mathbb{R})$ connection.

$$\omega = \int_{\Sigma} \text{Tr}(\delta A_1 \delta A_2)$$

Partition function

$$\mathcal{F}_g = \int_{\mathcal{M}_g} e^{\omega}$$

Related to tautological classes

$$\omega = \pi_*(\tau_2)$$

Contact terms

$$\left\langle \exp(y\omega) \right\rangle = \left\langle \exp \left(\sum_{k=2}^{\infty} \frac{(-1)^k y^{k-1}}{(k-1)!} \tau_k \right) \right\rangle$$

Spectral curve (Eynard)

$$y = \frac{\sinh(\sqrt{\xi}x)}{\sqrt{\xi}}$$

Jackiw-Teitelboim Gravity

Hyperbolic metrics

$$S = \int_{\Sigma} \phi \sqrt{g} (R + 1) + \int_{\partial\Sigma} \phi_b K$$

Constant curvature metrics

$$R = -1$$

Dual to Schwarzian theory on the boundary, eqv SYK model

$$\int d\tau \{F, \tau\}$$

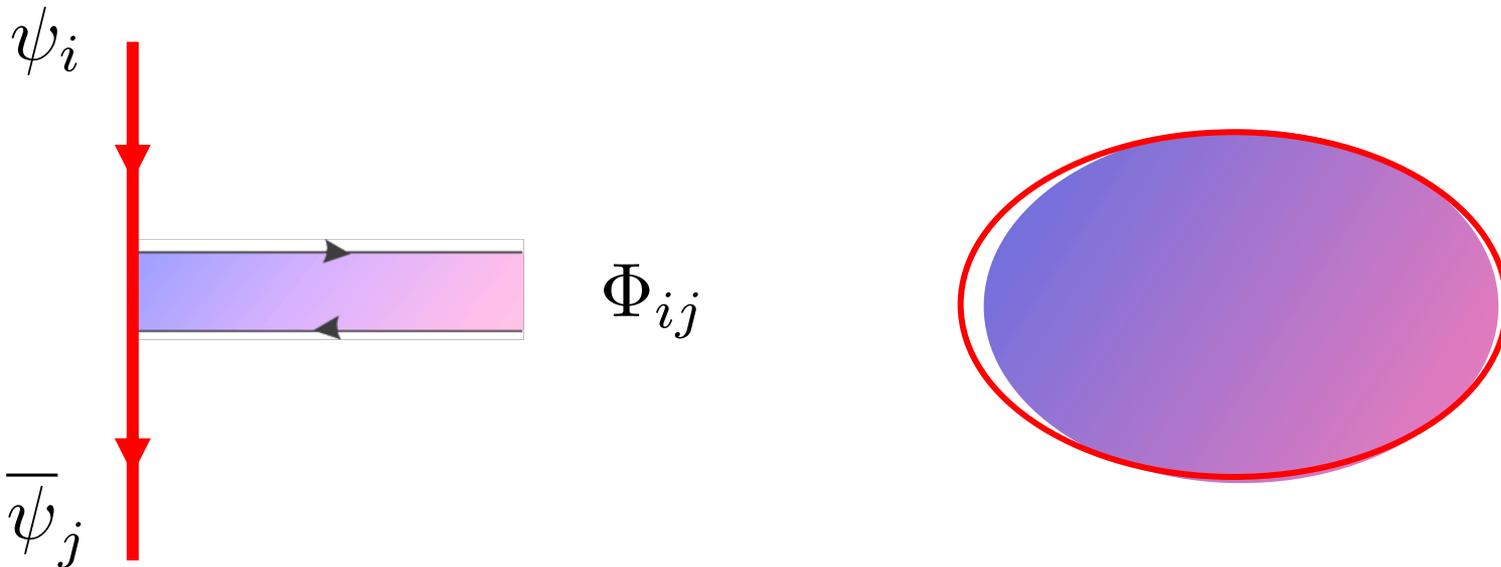
Witten-Stanford: $\rho(E) = \sinh \sqrt{E}$

Open Strings

Introduce vector degrees of freedom

$$\int d\Phi d\psi d\bar{\psi} e^{\text{Tr}W(\Phi) + \bar{\psi}(\Phi - z)\psi}$$

External quark line



Open Strings

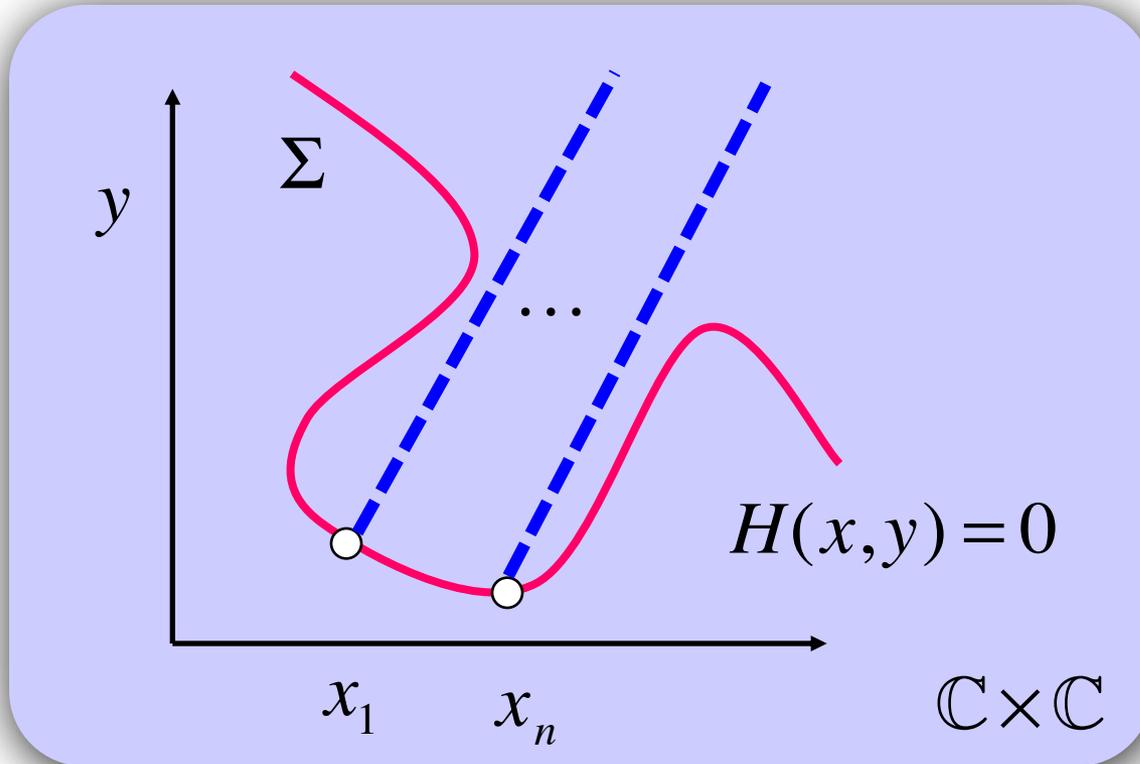
Integrate out vector-valued variables

$$\int d\Phi e^{\text{Tr}W(\Phi)} \det(\Phi - z)$$

Vertex operators = D-branes = fermions in chiral CFT

$$\psi(x) = e^{i\varphi(x)} = \det(x - \Phi)$$
$$\psi^*(x) = e^{-i\varphi(x)} = \det(x - \Phi)^{-1}$$

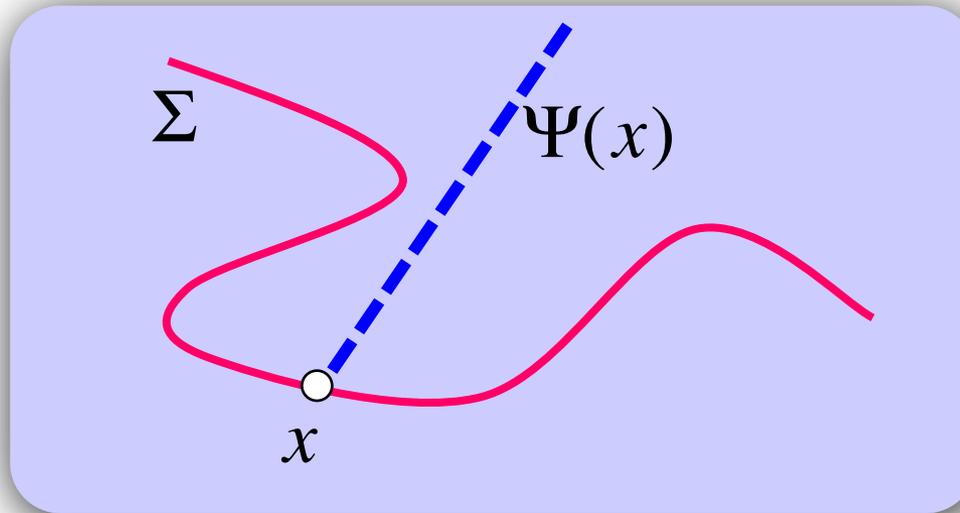
Topological D-Branes



Open string partition function

$$\Psi(x_1, \dots, x_n) = \langle \psi(x_1) \cdots \psi(x_n) \rangle = \left\langle \prod_i \det(\Phi - x_i) \right\rangle$$

Quantum Wave Function



Single brane wave function $\Psi(x)$

$$\hat{H}\Psi = 0, \quad \hat{H} = \hat{H}(\hat{x}, \hat{y})$$

Quantization of phase space

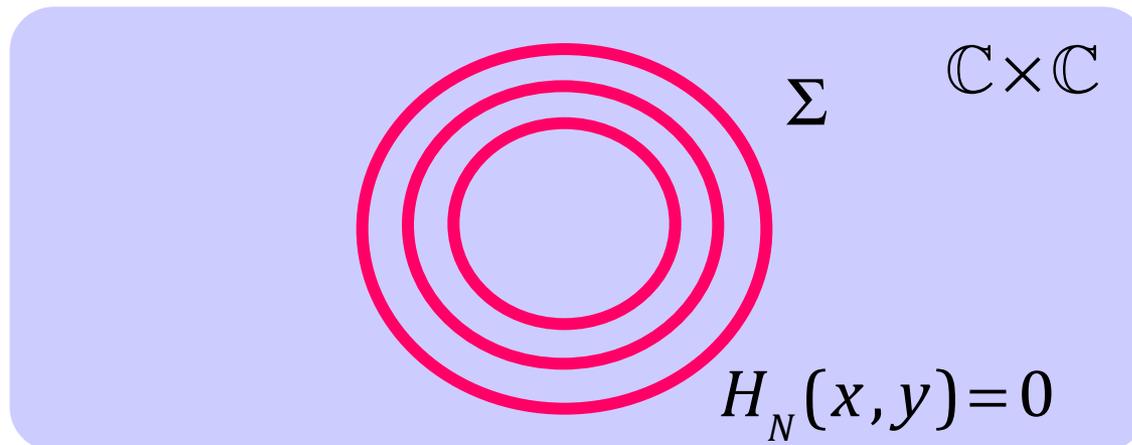
$$\hat{y} = -g_s \frac{\partial}{\partial \hat{x}}, \quad [\hat{x}, \hat{y}] = g_s$$

Gaussian Matrix Model

$$\Psi_N(x) = \left\langle \det(\Phi - x) \right\rangle_N = H_{N-1}(x) \cdot e^{-x^2/2}$$

Eigenfunctions of harmonic oscillator

$$\left(-g_s^2 \frac{\partial^2}{\partial x^2} + x^2 - g_s(2N-1) \right) \Psi_N(x) = 0$$



Topological Gravity

Airy equation

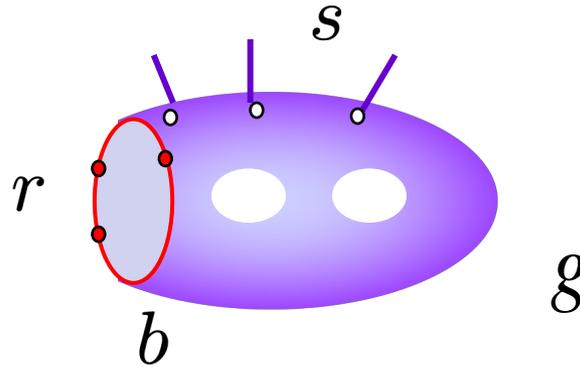
$$(\partial^2 - x)\Psi = 0$$

D-brane (open string) partition function

$$\Psi(z) = \int dp e^{izp + p^3/2}$$

Open Topological Strings

Surfaces with b holes, s bulk and r boundary punctures



Attempted definition

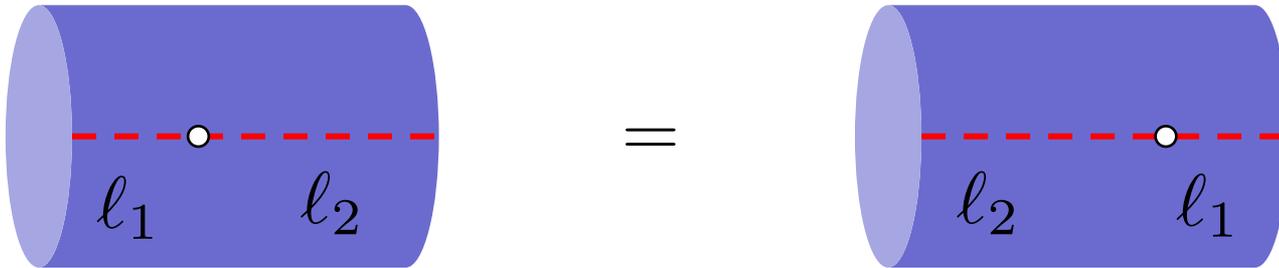
$$\langle \tau_{n_1} \cdots \tau_{n_s} \sigma^r \rangle = \int_{\mathcal{M}_{g,b,s,r}} c_1(L_1)^{n_1} \cdots c_1(L_s)^{n_s}$$

↑

Boundary puncture operator, dual to bdry cosm constant.

Open Topological String Anomaly

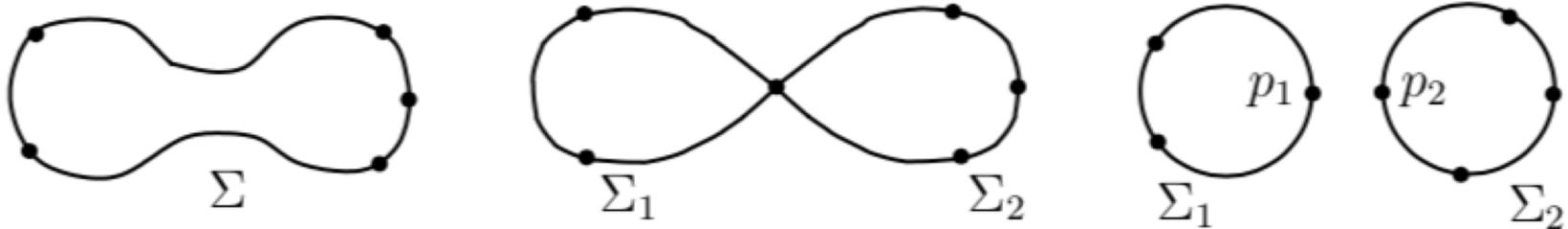
Moduli space is not orientable



$$dl_1 \wedge dl_2 = -dl_2 \wedge dl_1$$

Open Topological String Anomaly

Compactified moduli space has a boundary in (real) codim 1.
Degeneration of boundary:



Real dim $n - 3 \rightarrow \dim n - 4$

LG-model

Matter for $(p,1)$ model is topological gravity coupled to twisted $N=2$ SUSY minimal model (Landau-Ginzburg)

$$W(\phi) = \phi^p$$

In particular for “pure” topological gravity

$$W(\phi) = im\phi^2$$

After twisting W is spin one, so ϕ is spin $\frac{1}{2}$ (in general spin $1/p$).

Choice of spin structure

$$\phi \in S, \quad S^{\otimes 2} = K$$

Closed Surfaces

Spin structures are odd or even, measured by the number of fermion zero modes

$$w = \dim H^0(\Sigma, S) = 0, 1 \pmod{2}$$

Sum over spin structures for fixed genus g

$$\frac{1}{2} \sum_S (-1)^w = \frac{1}{2} \left(\frac{2^{2g} + 2^g}{2} - \frac{2^{2g} - 2^g}{2} \right) = 2^{g-1} = (\sqrt{2})^{2g-2}$$

Redefine the coupling constant

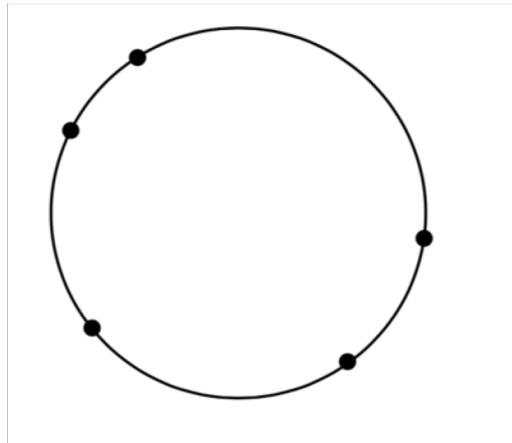
$$g_s \rightarrow \sqrt{2}g_s$$

Open Surfaces

Can we extend $(-1)^w$ to surfaces with boundaries?

Solved by mathematicians (Pandharipande, Solomon, Tessler)

Two choices of spin bundle on boundary. Include punctures that alternate between the two choices.



LG-model

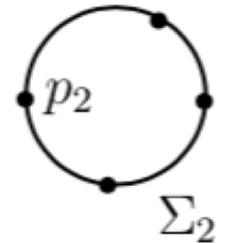
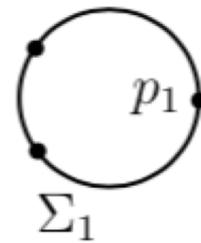
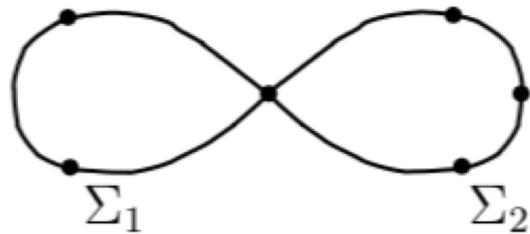
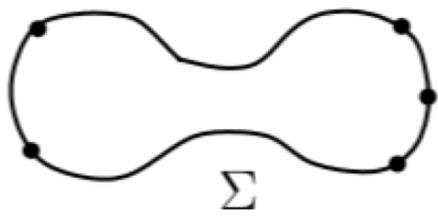
This prescription has a physical interpretation. In LG model there can be two choices of branes B, B' , depending on the orientation. Solutions of instanton equation

$$\bar{\partial}\phi = \bar{W}'$$

Quantize: only physical states in B - B' channel.
Super gauge group interpretation

$$\Phi = \begin{pmatrix} A & B \\ C & D \end{pmatrix}, \quad A, D \text{ even, } B, C \text{ odd.}$$

Anomalies



Matrix Models

Can derive the modified Virasoro constraints.

Open string partition function

$$\Psi(z) = \int dv \langle e^{v\sigma} \rangle$$

Satisfies

$$L_n \Psi(z) = \left(z^{n+1} \partial_z + \frac{1}{4} (n+1) z^n \right) \Psi(z)$$

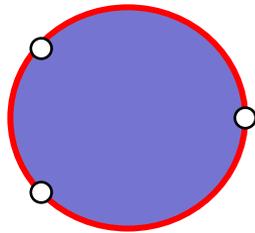
$$(\partial_z^2 - z + t_0) \Psi(z) = 0$$

Matrix Models

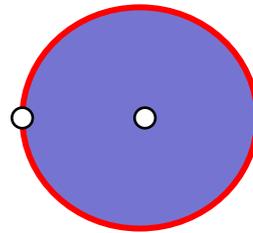
Consider the case of only punctures (bulk & boundary)

$$\langle e^{t_0 \tau_0 + v \sigma} \rangle = e^{(t_0 v + v^3/3)/g_s}$$

Airy function!



$$\langle \sigma^3 \rangle$$



$$\langle \sigma \tau_0 \rangle$$

Conclusions

The relation between matrix models and topological Gravity can be extended to open strings. Many subtle Effects to take care of.

Open problems:

- Extend Mirzakhani's work to unoriented strings. Problem: moduli space seems to be non-compact.
- Extend to topological supergravity. Moduli space of super-Riemann surfaces. Some preliminary results suggests that it's related to the (-2,1) model, or LG with singular potential

$$W(\phi) = \frac{1}{\phi^2}, \quad y^2 = \frac{1}{x}$$