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 (Prof. E. Frey)

Problem set 6

Problem 6.1 *effective free energy of a fluid*

In terms of the canonical ensemble, a fluid is characterized by a temperature T as conjugate variable to the energy, as well as the two macro-variables volume V , and particle number N .

Here we want to achieve a proper description of the fluid that includes small fluctuations of the temperature $T(\vec{r}) = T + \delta T(\vec{r})$ and the local particle densities $n(\vec{r}) = n + \delta n(\vec{r})$ close to their equilibrium values. Up to quadratic order an effective hamiltonian is given by

$$\mathcal{H}_{\text{eff}} = \frac{1}{2} \int d^d \vec{r} \left[\frac{C_V}{VT} \delta T(\vec{r})^2 + \frac{1}{n^2 \kappa_T} \delta n(\vec{r})^2 \right],$$

where C_V denotes the heat capacity at constant volume, and κ_T is the isothermal compressibility.

1. Introduce discrete Fourier modes in the finite box $V = L^d$,

$$\delta n_{\vec{k}} = \frac{1}{V} \int_V d^d \vec{r} e^{-i\vec{k} \cdot \vec{x}} n(\vec{r}), \quad \delta n(\vec{r}) = \sum_{\vec{k}} e^{i\vec{k} \cdot \vec{x}} n_{\vec{k}}$$

and similarly for $\delta T(\vec{r})$ to obtain a representation of \mathcal{H}_{eff} in Fourier space.

2. Calculate thermal averages of the density, temperature modes, e.g. $\langle n_{\vec{k}} n_{\vec{q}} \rangle$ using the equipartition theorem.
3. Use thermodynamic arguments to argue that the effective hamiltonian gives a faithful representation of the macroscopic properties of a fluid. In particular, corroborate that \mathcal{H}_{eff} should not contain mixed terms $\delta T(\vec{r}) \delta n(\vec{r})$.

Problem 6.2 *tricritical point*

Consider the partition sum for a single scalar order parameter field $m(\vec{x})$,

$$Z = \int \mathcal{D}[m(\vec{x})] \exp(-\beta \mathcal{H}[m(\vec{x})]),$$

for the Ginzburg-Landau functional

$$\beta \mathcal{H}[m(\vec{x})] = \int d\vec{x} \left\{ \frac{a}{2} m(\vec{x})^2 + \frac{c}{2} (\vec{\nabla} m(\vec{x}))^2 + \frac{u}{4} m(\vec{x})^4 + \frac{v}{6} m(\vec{x})^6 \right\},$$

where now the phenomenological parameters a, u may take either sign, whereas stability is guaranteed by the inequalities $c > 0$ and $v > 0$.

1. Perform a functional derivative of $\beta \mathcal{H}$ to derive an equation for the configuration $m(\vec{x})$ that minimizes $\beta \mathcal{H}$.

2. Consider now only homogeneous solutions $m(x) = \text{const.}$. Make a sketch of the effective free energy and show that there may be five extremal points. Discuss their respective stability and distinguish a continuous and discontinuous phase transition scenario. Make a sketch of the phase diagram in the a - v plane and indicate the nature of the phase boundaries.
3. The *tricritical* point $(a, u) = (0, 0)$ is a special point in the phase diagram. Extend the Ginzburg-Landau functional by a coupling to some space-varying external field $h(\vec{x})$,

$$\beta\mathcal{H}[m(\vec{x})] = \int d\vec{x} \left\{ \frac{a}{2} m(\vec{x})^2 + \frac{c}{2} (\vec{\nabla} m(\vec{x}))^2 + \frac{u}{4} m(\vec{x})^4 + \frac{v}{6} m(\vec{x})^6 - h(\vec{x}) m(\vec{x}) \right\}.$$

Assuming a linear temperature dependence for $a = a'\tau$, and $u = u'\tau$ with the reduced temperature $t = (T - T_c)/T_c$, discuss the critical behavior close to the tricritical point. In particular, determine the critical exponents

- (a) for the low-temperature phase behavior of the order parameter $m \sim (-\tau)^{\beta_t}$,
- (b) the zero-field susceptibility $\chi \sim |\tau|^{-\gamma_t}$,
- (c) and the divergence of the correlation length $\xi \sim |\tau|^{-\nu_t}$.
- (d) Furthermore show that at the tricritical point the order parameter depends singularly on the field $m \sim h^{1/\delta_t}$.