On Mirror Symmetry for Calabi-Yau Fourfolds with Three-Form Cohomology

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arXiv: 1512.04859 (T. Grimm, SG)

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Mirror symmetry of CY_4

F-theory effective action on elliptically fibered CY_4 :

 $\mathcal{N}=1$ supergravity in (3+1) dimensions

$$S^{(4)} = \int_{\mathcal{M}_{3,1}} \frac{1}{2} R^{(4)} * 1 - \mathcal{K}_{IJ}^{F} \mathcal{D} \mathcal{M}^{I} \wedge * \mathcal{D} \bar{\mathcal{M}}^{J} \\ - \frac{1}{2} Re(f)_{\Lambda \Sigma} \mathcal{F}^{\Lambda} \wedge * \mathcal{F}^{\Sigma} - \frac{1}{2} Im(f)_{\Lambda \Sigma} \mathcal{F}^{\Lambda} \wedge \mathcal{F}^{\Sigma} + V * 1$$

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For gauge-theory on seven-branes \Rightarrow not relevant

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What can we say about the gauge-coupling $f_{\Lambda\Sigma}(M)$ in terms of CY_4 ? For gauge-theory on seven-branes \Rightarrow not relevant

For bulk gauge-theory: \Rightarrow Emerge from non-trivial three-forms of CY_4

Goal: Study three-forms of $CY_4 \Rightarrow$ Consider toy-example: IIA on CY_4 !

Type IIA Supergravity

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- 2 Properties of CY_4

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- Summary and Outlook

Three parts:

$$S_{IJA}^{(10)} = S_{NS} + S_{kin} + S_{CS}$$

 $S_{IIA}^{(10)}=S_{NS}+S_{kin}+S_{CS}$

Neveu-Schwarz

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Ramond-part (gauge-kinetic)

$$S_{\textit{kin}} = -\frac{1}{4}\int_{\mathcal{M}_{9,1}}F_2 \wedge \ast F_2 + \tilde{F}_4 \wedge \ast \tilde{F}_4$$

with
$$F_2 = dC_1$$
, $\tilde{F}_4 = F_4 - C_1 \wedge H_3$

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Characteristic feature: SU(4)-holonomy

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Covariantly constant Weyl-spinor:

 $abla \eta = 0, \quad \gamma_9 \eta = \eta$

Hodge diamond:



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Calabi-Yau fourfold compactifications due to [Haack,Louis,...]

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Kähler deformations

$$g_{i\overline{j}} + \delta g_{i\overline{j}} = -iJ_{i\overline{j}} = -i\sum_{A=1}^{h^{1,1}} v^A(x)\omega_{Ai\overline{j}}, \quad \omega_A \in H^2(Y_4,\mathbb{Z})$$

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 $h^{1,1}$ complexified Kähler moduli:

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$$t^A = b^A + i v^A$$

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Without further moduli:
$$\mathcal{M}_{moduli} = \mathcal{M}_k \times \mathcal{M}_{cs}$$

$$e^{-2\phi_{IIA}^{(10)}(x)} = rac{1}{\mathcal{V}(x)}e^{-2\phi_{IIA}^{(2)}(x)}, \quad \mathcal{V} = rac{1}{4!}\int_{Y_4}J^4$$

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 \Rightarrow Choose three-form basis Ψ^{I} depending on complex structure moduli!

$$\Psi'(z,\overline{z}) = \frac{1}{2} \operatorname{Re}(f)^{lm} (\alpha_m - i\overline{f}_{mk}\beta^k) \in H^{1,2}(Y_4)$$

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Properties:

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Advantage: $C_{AI}^{k} = \int \omega_{A} \wedge \alpha_{I} \wedge \beta^{k}$ topological intersection numbers

$$\Rightarrow \quad \int_{Y_4} \Psi' \wedge * \overline{\Psi}^k = -\frac{1}{2} Re(f)^{lm} C^k_{Am} v^A$$

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$$\Rightarrow \int_{Y_4} \Psi' \wedge * \overline{\Psi}^k = -\frac{1}{2} Re(f)^{lm} C^k_{Am} v^A \qquad \begin{array}{c} \text{Goal of our work:} \\ \text{calculate } f! \end{array}$$

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Dimensional Reduction on $\mathbb{M}_{1,1} \times Y_4$

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with kinetic potential $K = \log \mathcal{V} - \log \int_{Y_4} \Omega \wedge \overline{\Omega} + e^{2\phi_{IIA}^{(2)}} S$

$$\mathcal{S} = Re(N)_I \Big(\int_{Y_4} \Psi^I \wedge * \overline{\Psi}^m \Big) Re(N)_m \quad \Rightarrow \quad Im(N)_I ext{ axionic!}$$

 \Rightarrow Interactions between chiral and twisted-chiral multiplets!

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New: Three-form moduli are self-mirror, but change representations!

 \Rightarrow Representation change via Legendre transformation of kinetic potential \Rightarrow Explicit dependence of three-form metric

$$H^{lk} = \int_{Y_4} \Psi^l \wedge * \overline{\Psi}^k = -rac{1}{2} Re(f)^{lm} C^k_{Am} v^A = (H^{-1}_{\mathsf{mirror}})_{lk}$$

on Kähler moduli v^A in both pictures known \Rightarrow Gauge coupling f(z) linear in complex structure moduli z!

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Summary

Introducing the ansatz

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we can include the massless-modes arising from non-trivial massless three-form modes into dimensional reductions on Calabi-Yau fourfolds.

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we can include the massless-modes arising from non-trivial massless three-form modes into dimensional reductions on Calabi-Yau fourfolds.

Using type IIA supergravity and $\mathcal{N} = (2, 2)$ supersymmetry, we obtain a very simple mirror map. This enables us to calculate $f_{lm}(z)$ at large complex structure to be

$$\hat{f}_{lm}(z) = z^{K} \hat{C}_{Kl}^{m} \qquad \qquad \hat{C}_{Kl}^{m} = \int_{\hat{Y}_{4}} \hat{\omega}_{K} \wedge \hat{\alpha}_{l} \wedge \hat{\beta}^{m}$$

(Hatted objects are defined on the mirror fourfold \hat{Y}_4)

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 - \Rightarrow three-forms induce $U(1)^{h_{2,1}}$ -gauge theory in the effective supergravity action

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and we calculated $f_{\Lambda\Sigma}(z)$ and K^F at large complex structure!

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• Continued in [Corvilain, Grimm, Regalado '16].

Thank you for your attention!

Questions?

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 - \Rightarrow done! hypersurfaces in toric varieties
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Distant future:

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Distant future:

- Proper geometrical interpretation of three-form moduli
 - \Rightarrow probably moduli of quasi-coherent sheaves or branes
 - \Rightarrow generalize to open/homological mirror-symmetry concepts

- Construct examples via toric geometry
 - \Rightarrow done! hypersurfaces in toric varieties
 - \Rightarrow three-forms arise from blow-ups of singular Riemann surfaces
 - \Rightarrow calculate three-form periods as regular periods of these surfaces
- Apply ansatz to M/F-theory
 - \Rightarrow phenomenological applications: axions for inflation
 - \Rightarrow include open string moduli and fluxes
 - \Rightarrow new information about brane dynamics

Distant future:

- Proper geometrical interpretation of three-form moduli
 - \Rightarrow probably moduli of quasi-coherent sheaves or branes
 - \Rightarrow generalize to open/homological mirror-symmetry concepts
- Include three-form moduli in topological string/CFT framework