SUPERSYMMETRIC FLAVOUR MODELS

with
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arXiv:1308.1090 [hep-ph]

- M. Goodsell, L. Heurtier and P. Tziveloglou, in progress .

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Outline

- Supersymmetry and naturalness
- Flavor and inverted hierarchy/natural SUSY

- U(1) and U(2) models
- U(1) x (discrete subgroup of) SU(2)

- Dirac gauginos and flavor models

- UV softness and flavor suppression with Dirac gauginos
- Mass scales
- Flavor U(1) models, U(1) x SU(2)

- Conclusions

Supersymmetry and Naturalness

SUSY is still the simplest and most elegant solution to the hierarchy problem.

$$\delta m_h^2 \approx \frac{3g^2 m_t^4}{8\pi^2 m_W^2}\left[\ln\left(\frac{M_{\rm SUSY}^2}{m_t^2}\right) + \frac{X_t^2}{M_{\rm SUSY}^2}\left(1 - \frac{X_t^2}{12M_{\rm SUSY}^2}\right)\right]$$

where and M_{SUSY} (A_t) denotes the average stop mass (mass mixing in the stop sector).

Electroweak scale natural for light higgsinos,gluinos,stops and L-handed sbottom:

$$m_Z^2 = -2(m_{H_u}^2 + |\mu|^2) + \ldots$$

$$\delta m_{H_u}^2 \approx -\frac{3y_t^2 m_{\tilde{t}}^2}{4\pi^2}(1 + a^2/2)\log\frac{\Lambda}{m_{\tilde{t}}}$$

$$\delta m_{\tilde{t}}^2 = \frac{8\alpha_s}{3\pi} M_3^2 \log\frac{\Lambda}{M_3}$$

(More) Natural SUSY models:

- Natural SUSY/inverted hierarchy/split families :

light stops,gluinos,higgsinos (TeV)

heavier 1,2 generations (10-15 TeV)

- Extended scalar and/gauge sector (ex: NMSSM)
- RPV models (ex. baryonic RPV, operators UDD)
- Dirac gauginos
- Spectrum more degenerate/decays stealthy…

(Less) Natural SUSY theories :

- Mini-split/Spread SUSY models
- Split SUSY models: $m_{\rm scalars} >> m_{\rm fermions}$
- High-scale SUSY

SUSY comments from LHC searches and SM scalar mass :

- LHC direct SUSY searches and Higgs mass set new limits on superpartner masses for simple (simplified) SUSY models $m_{gluinos}, m_{squarks} \geq 1.5\ TeV$

Popular models: mSUGRA, CMSSM, minimal gauge mediation with TeV superpartner masses have difficulties in accomodating the data in a natural way .

- However, from a UV viewpoint (supergravity, string theory), popular models are unnatural.

It is important to theoretically analyze and experimentally search for non-minimal SUSY models.

Inverted hierarchy/Natural SUSY

An old scenario which became popular recently because of LHC constraints:

- third generations squarks and gauginos in the TeV range (light stops).
- First two generation scalars much heavier (10-15 TeV). They affect little the tuning of the electroweak scale.

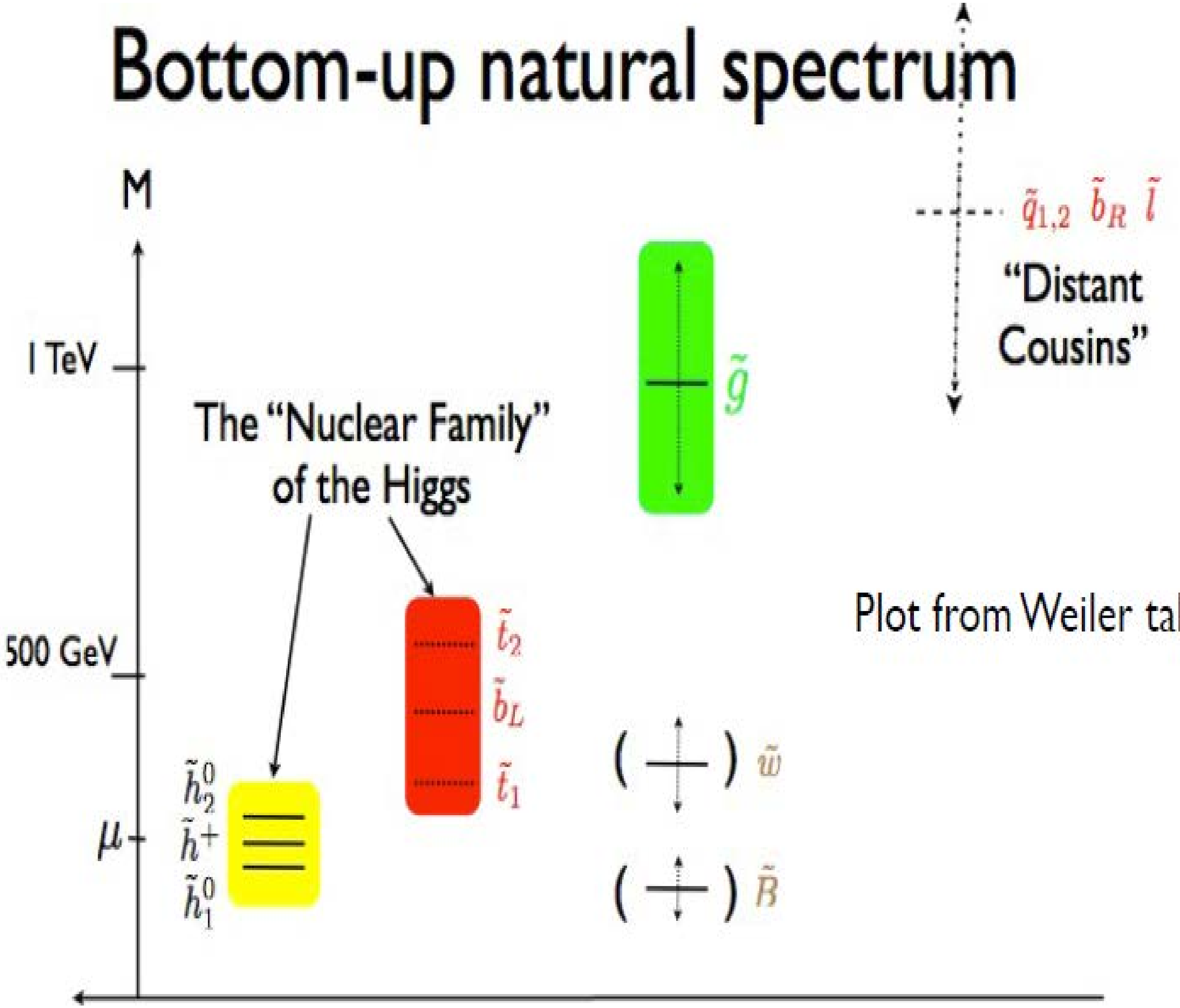
This is natural in flavor models and holographic constructions.

1) Simplest flavor model: U(1) gauged flavor symmetry (Froggatt-Nielsen,79). Quark mass matrices given by

$$h_{ij}^U \sim \epsilon^{q_i+u_j+h_u} \quad , \quad h_{ij}^D \sim \epsilon^{q_i+d_j+h_d} \ ,$$

where typically $\epsilon = \frac{\langle \Phi \rangle}{M} \sim \lambda = 0.22$ and q_i are charges of left-handed quarks, etc.

Bottom-up natural spectrum



Plot from Weiler talk on natural susy

Quarks masses and mixings are given by ($q_{13} = q_1 - q_3$,etc)

$$\frac{m_u}{m_t} \sim \epsilon^{q_{13}+u_{13}} \, , \quad \frac{m_c}{m_t} \sim \epsilon^{q_{23}+u_{23}} \, , \quad \frac{m_d}{m_b} \sim \epsilon^{q_{13}+d_{13}} \, , \quad \frac{m_s}{m_b} \sim \epsilon^{q_{23}+d_{23}}$$

$$\sin\theta_{12} \sim \epsilon^{q_{12}} \quad , \quad \sin\theta_{13} \sim \epsilon^{q_{13}} \quad , \quad \sin\theta_{23} \sim \epsilon^{q_{23}} \, .$$

Good fit to to data ⇨ larger charges for the lighter generations

$$q_1 > q_2 > q_3 \quad , \quad u_1 > u_2 > u_3 \quad , \quad d_1 > d_2 > d_3$$

$$m_t \sim 1 \qquad m_c \sim \epsilon^4 \qquad m_u \sim \epsilon^8$$

$$m_b \sim \epsilon^3 \qquad m_s \sim \epsilon^{5\div 6} \qquad m_d \sim \epsilon^{7\div 8}$$

$$m_\tau \sim \epsilon^3 \qquad m_\mu \sim \epsilon^5 \qquad m_e \sim \epsilon^9$$

$$V_{us} \sim \epsilon \qquad V_{ub} \sim \epsilon^3 \qquad V_{cb} \sim \epsilon^2$$

Gauge anomalies → constraints on the charges, Green-Schwarz mechanism, anomalous U(1)

$$K \sim \frac{X^\dagger X}{\Lambda_S^2} \left(\frac{\phi}{\Lambda_F} \right)^{|q_i - q_j|} Q_i^\dagger Q_j$$

→ F-term contributions to scalar masses.

There are also D-term contributions, so

scalar masses are of the form

$$m_{ij}^2 \ = \ X_i \langle D \rangle \ + \ a_{ij} \langle F \rangle$$

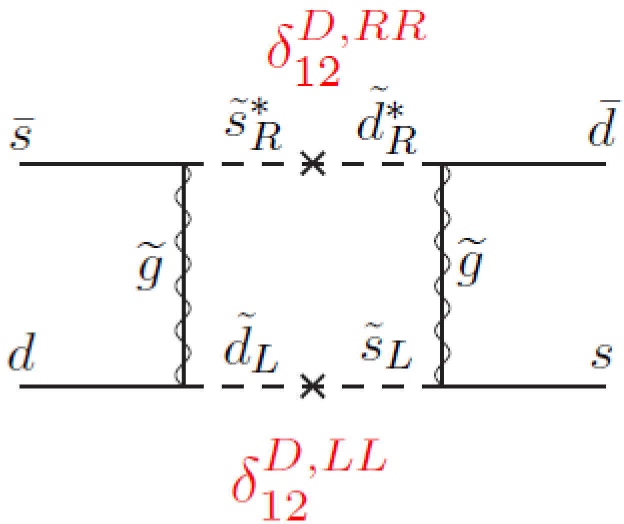
If $D >> F$ then an inverted hierarchy is generated.

This can be realized in explicit models (E.D.,Pokorski,Savoy; Binetruy,E.D.; Dvali,Pomarol,94-96)

Obs: 1-2 generations cannot be too heavy → tachyonic stops (Pomarol,Tommasini;Arkani-Hamed,Murayama)

.

Nowdays, FCNC constrain seriously these models; need some degeneracy 1,2 generations.



$$\sim \frac{1}{m^2} \quad \delta_{12}^{D,LL} \; \delta_{12}^{D,RR} \quad \sim \frac{1}{m^2} \quad \frac{m_d}{m_s}$$

⇨ $q_1 = q_2 \; , q_3 = 0$ if not m > 100 TeV or so.

But then m_{12}^2 squark mass not protected by the U(1) symmetry

There is a challenge to explain simultaneously fermion masses and FCNC within one flavour theory !

Operator	Bounds on Λ in TeV ($c_{ij}=1$)		Bounds on c_{ij} ($\Lambda = 1$ TeV)		Observables
	Re	Im	Re	Im	
$(\bar{s}_L \gamma^\mu d_L)^2$	9.8×10^2	1.6×10^4	9.0×10^{-7}	3.4×10^{-9}	Δm_K; ϵ_K
$(\bar{s}_R\, d_L)(\bar{s}_L d_R)$	1.8×10^4	3.2×10^5	6.9×10^{-9}	2.6×10^{-11}	Δm_K; ϵ_K
$(\bar{c}_L \gamma^\mu u_L)^2$	1.2×10^3	2.9×10^3	5.6×10^{-7}	1.0×10^{-7}	Δm_D; $\lvert q/p\rvert, \phi_D$
$(\bar{c}_R\, u_L)(\bar{c}_L u_R)$	6.2×10^3	1.5×10^4	5.7×10^{-8}	1.1×10^{-8}	Δm_D; $\lvert q/p\rvert, \phi_D$
$(\bar{b}_L \gamma^\mu d_L)^2$	5.1×10^2	9.3×10^2	3.3×10^{-6}	1.0×10^{-6}	Δm_{B_d}; $S_{\psi K_S}$
$(\bar{b}_R\, d_L)(\bar{b}_L d_R)$	1.9×10^3	3.6×10^3	5.6×10^{-7}	1.7×10^{-7}	Δm_{B_d}; $S_{\psi K_S}$
$(\bar{b}_L \gamma^\mu s_L)^2$	1.1×10^2		7.6×10^{-5}		Δm_{B_s}
$(\bar{b}_R\, s_L)(\bar{b}_L s_R)$	3.7×10^2		1.3×10^{-5}		Δm_{B_s}

TABLE I: Bounds on representative dimension-six $\Delta F = 2$ operators. Bounds on Λ are quoted assuming an effective coupling $1/\Lambda^2$, or, alternatively, the bounds on the respective c_{ij}'s assuming $\Lambda = 1$ TeV. Observables related to CPV are separated from the CP conserving ones with semicolons. In the B_s system we only quote a bound on the modulo of the NP amplitude derived from Δm_{B_s} (see text). For the definition of the CPV observables in the D system see Ref. [15].

2) FCNC constraints are better enforced by non-abelian symmetries.

A popular example: U(2) = SU(2) x U(1) flavor symmetry (Pomarol,Tommasini; Barbieri,Dvali,Hall…)

- 1st,2nd generations : U(2) doublets, scalars degenerate
- 3rd generation: singlet

Here, FCNC are largely suppressed.

$$\Delta C_1 \sim \frac{\alpha_s^2}{m_{\tilde g}^2} \left[(Z_D^L)^*_{13} (Z_D^L)_{23} \right]^2 \left[f_4(x_1,x_1) - 2 f_4(x_1,x_3) + f_4(x_3,x_3) \right]$$

$\underbrace{(Z_D^L)^*_{13}(Z_D^L)_{23}}$: $\hat\delta_{12}^{D,LL}$ ⟵ known from fermion sector

$$\sim V_{cb}^2 \sqrt{m_d/m_s}$$

where $x_i = \frac{m_i^2}{M_{\tilde g}^2}$ and f_4 are loop functions.

However, there are two problems :

- One with the CKM elements:

$$|V_{td}/V_{ts}| = \sqrt{m_d/m_s}\ [1+\mathcal{O}(\epsilon^2)]$$

0.22 ± 0.01 $\quad$ 0.22 ± 0.02

$\mathcal{O}(10^{-3})$

$$|V_{ub}/V_{cb}| = \sqrt{m_u/m_c}\ [1+\mathcal{O}(\epsilon^2)]$$

0.085 ± 0.004 $\quad$ 0.046 ± 0.008

- Another possible problem : $\tan\beta$ typically large. Then the minimal natural SUSY spectrum with heavy $\tilde{b}_R$ has difficulties with RG running from GUT to EW scale

Possible to combine abelian+non-abelian flavor symmetries in a constructive way: U(1) x D'_n, where D'_n is a discrete non-abelian subgroup of SU(2) (DGPZ)

$U(1)$ — protect ϵ_K → $SU(2) \times U(1)$ ← correct V_{ub}/V_{cb} — $U(2)$

Split spectrum from U(1) D-term

$\tilde{m}_D$ — $\sim 10\,\mathrm{TeV}$ $\quad \tilde{q}_{1,2}, \tilde{b}_R, \tilde{\tau}_L$

$\tilde{m}_F$ — $\sim 1\,\mathrm{TeV}$ $\quad \tilde{t}_L, \tilde{t}_R, \tilde{b}_L, \tilde{\tau}_R \quad \mu, B_\mu, \tilde{m}^2_{H_u}, \tilde{m}^2_{H_d}, M_i, A$

	10_a	10_3	$\overline{5}_a$	$\overline{5}_3$	H_u	H_d	ϕ^a	χ
$SU(2)$	2	1	2	1	1	1	$\overline{2}$	1
$U(1)$	X_{10}	0	$X_{\overline{5}}$	X_3	0	0	X_ϕ	-1

Table 1: Flavor group representations of the model.

$$\langle \phi^a \rangle = \epsilon_\phi \Lambda \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad , \qquad \langle \chi \rangle = \epsilon_\chi \Lambda$$

Model	ϵ_ϕ	ϵ_χ	$\tan\beta$	X_ϕ	X_{10}	X_5	X_3
A	0.02	0.02	5	-1	1	1	1
B	0.1	0.2	5	-2	3	3	2
B′	0.1	0.2	20	-2	3	2	1
C	0.2	0.1	50	-1	2	1	0

Table 2: Possible choices of parameters compatible with the fit to fermion masses and mixings.

The Yukawa matrices are given by

$$Y_u = \begin{pmatrix} 0 & h_{12}^u \epsilon_u' & 0 \\ -h_{12}^u \epsilon_u' & h_{22}^u \epsilon_u^2 & h_{23}^u \epsilon_u \\ 0 & h_{32}^u \epsilon_u & h_{33}^u \end{pmatrix} ,$$

$$Y_d = \begin{pmatrix} 0 & h_{12}^d \epsilon_u' \epsilon_d / \epsilon_u & 0 \\ -h_{12}^d \epsilon_u' \epsilon_d / \epsilon_u & h_{22}^d \epsilon_u \epsilon_d & h_{23}^d \epsilon_3 \epsilon_u \\ 0 & h_{32}^d \epsilon_d & h_{33}^d \epsilon_3 \end{pmatrix} ,$$

with

$$\epsilon_u \equiv \epsilon_\phi \epsilon_\chi^{X_{10}+X_\phi}, \qquad \epsilon_d \equiv \epsilon_\phi \epsilon_\chi^{X_{\bar{5}}+X_\phi}, \qquad \epsilon_u' \equiv \epsilon_\chi^{2X_{10}}, \qquad \epsilon_3 \equiv \epsilon_\chi^{X_3}.$$

We find

$$\mathrm{Im}\,\Delta C_4 \approx \frac{2}{3}\alpha_s^2 \frac{m_d}{m_s} |V_{23}^d|^2 s_d^2 \sin 2\tilde{\alpha}_{12} \left(\tilde{m}_{d_R}^2 - \tilde{m}_{b_R}^2\right) \frac{\log\left(\frac{\tilde{m}_{d_R}}{m_{\tilde{g}}}\right) + \frac{1}{4}}{(\tilde{m}_{d_R})^4}$$

$$\approx 1.6 \times 10^{-8} \left(\frac{|V_{23}^d|}{0.04}\right)^2 \left(\frac{s_d^2}{0.2}\right) \left(\frac{\sin \alpha_{12}}{0.5}\right) \left(\tilde{m}_{d_R}^2 - \tilde{m}_{b_R}^2\right) \frac{\log\left(\frac{\tilde{m}_{d_R}}{m_{\tilde{g}}}\right) + \frac{1}{4}}{(\tilde{m}_{d_R})^4}$$

where $t_d \equiv \tan\theta_d \equiv \dfrac{|h_{32}^d|\epsilon_d}{|h_{33}^d|\epsilon_3}$

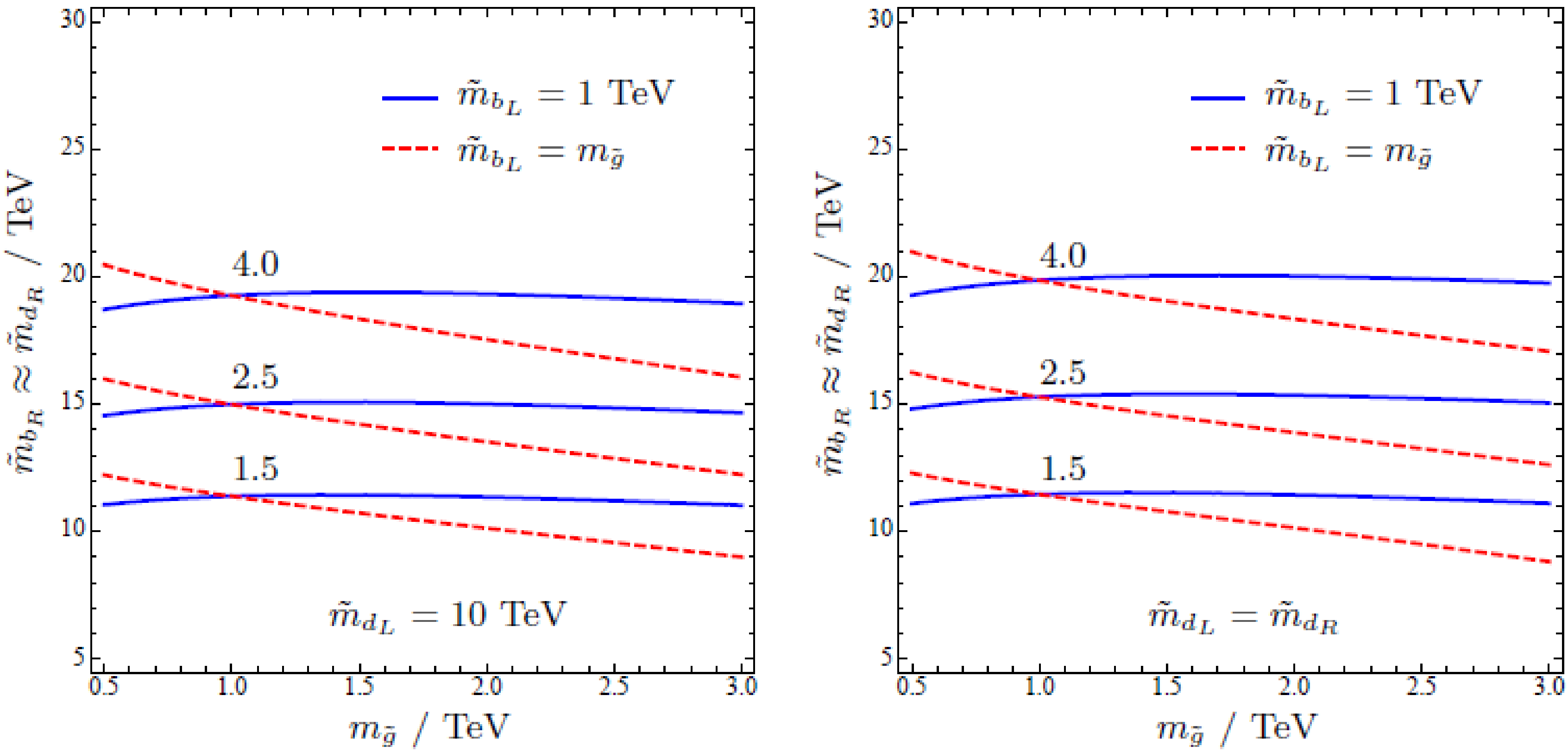


Figure 1: Bounds on the masses of the gluino and the approximately degenerate right handed down squark sector for various choices of the parameters. The region below each line is excluded. The three lines correspond to different choices of the dominant 3-1 splitting, namely $\tilde{m}^2_{d_R} - \tilde{m}^2_{b_R} = (1.5,\ 2.5,\ 4.0\ \text{TeV})^2$. The remaining parameters are chosen as $|V^d_{23}| = 0.04$, $\sin(\alpha_{12}) = 0.5$ and $s^2_d = 0.2$. The decoupling of the gluino occurs outside the displayed range of the gluino mass.

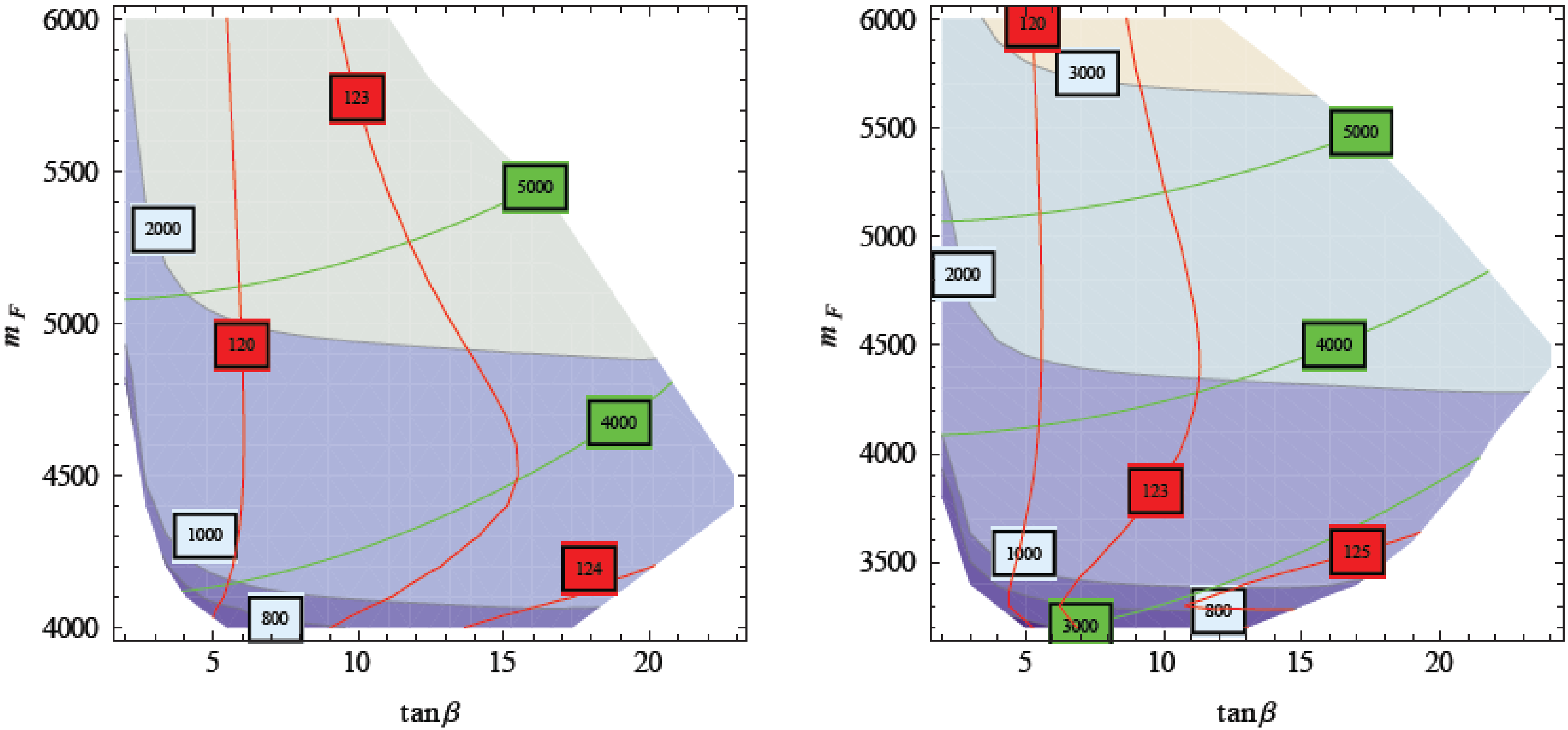
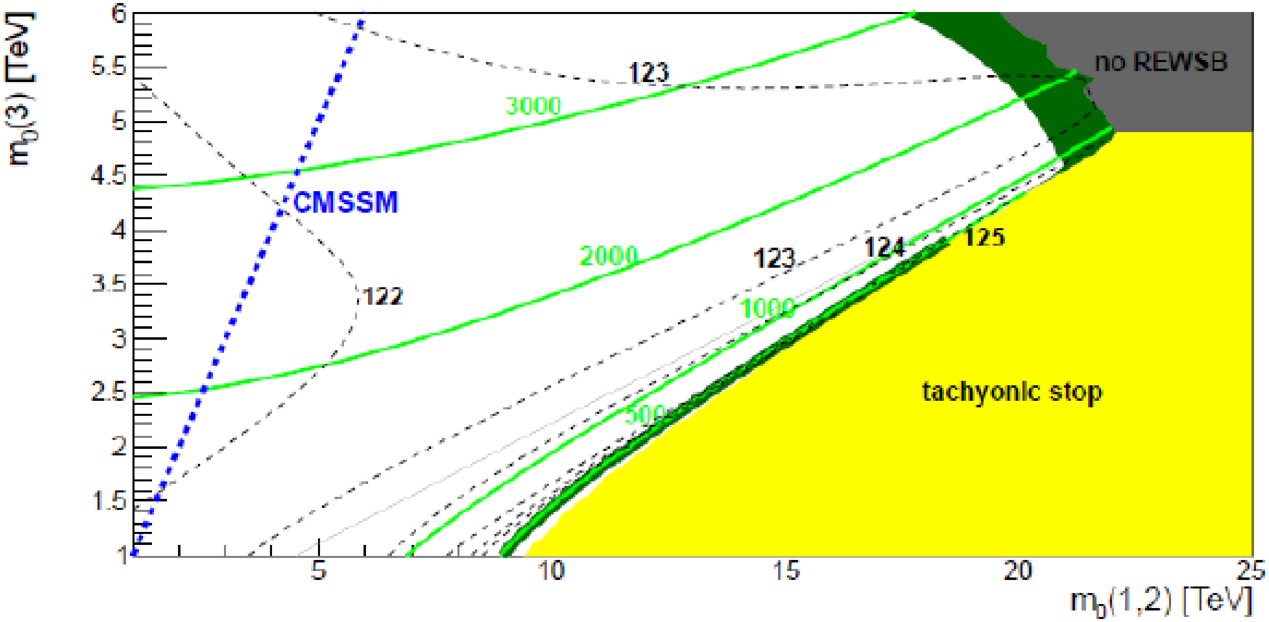


Figure 3: Parameter region in the $\tilde{m}_F/\tan\beta$ plane for fixed $\tilde{m}_D = 15\,\text{TeV}$ and $M_{1/2} = 0.6\,\text{TeV}$ (left panel) and $M_{1/2} = 1.0\,\text{TeV}$ (right panel). The contour lines correspond to the masses of $\tilde{t}_1$ (blue), $\tilde{\tau}_1$ (green) and h^0 (red).

(Courtesy of M. Badziak)

Large stop mixing can be generated from RG running (M. Badziak et al, 2012; Brummer et al, 2012.)



Inverted hierarchy example. Higgs mass (black dashed),stop mass (solid green) for $\mu > 0, \tan\beta = 10, M_{1/2} = 1, A_0 = -2$ (TeV). Yellow "tachyonic stop" and grey "no REWSB" ($\mu^2 < 0$) regions are excluded. Dark green region: $\Omega_{\mathrm{DM}} h^2 < 0.1288$.

Some issues model building: ::

- Discrete subgroups★ D'_n of U(2) avoid Goldstone bosons
- Simplest working example: D'_n with 12 elements generated by 2 generators with

$$A^6 = 1 \qquad B^4 = 1 \qquad ABA = B$$

On 2-dim. representations, they act as

$$\mathbf{2}_1 : \begin{pmatrix} e^{\pi i/3} & 0 \\ 0 & e^{-\pi i/3} \end{pmatrix}, \quad \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix}$$

$$\mathbf{2}_2 : \begin{pmatrix} e^{2\pi i/3} & 0 \\ 0 & e^{-2\pi i/3} \end{pmatrix}, \quad \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

- Operators breaking SU(2), invariant under D'_n appear usually at higher order in the lagrangian.

★Recent progress in string realization: Nilles et al, Camara et al…

Dirac gauginos and flavor models

SUSY with Dirac gauginos = MSSM + chiral adjoints χ

- Matter sector has N=1 spectrum
- Gauge multiplets are in N=2 multiplets
- There are Dirac masses for gauginos/gluinos

$$\int d^2\theta \ W'_\alpha Tr(W^\alpha \chi) \rightarrow m_D \ Tr(\lambda \Psi_\chi)$$

where $W'_\alpha = \theta_\alpha m_D$ is a fermionic spurion.

- UV softness and flavor suppression with Dirac gauginos

Squark/slepton masses generated at one-loop are UV finite

$$m_{\tilde{f}}^2 = \sum_a \frac{C_a(f)\alpha_a}{\pi}(m_D^a)^2 \ln \frac{\tilde{m}_i^2}{(m_D^a)^2})$$

where $\tilde{m}_i^2$ is the SUSY breaking mass of $Re\ O_\chi$

- Due to UV finiteness, quantum corrections to squarks are smaller than in MSSM. It is possible parametrically than $m_{\tilde{q}} < m_D$ at low energy.

- In MSSM, corrections to the Higgs soft terms are enhanced by a large factor $\ln \frac{\Lambda}{M_3}$. In purely Dirac MSSM case, the correction is softer

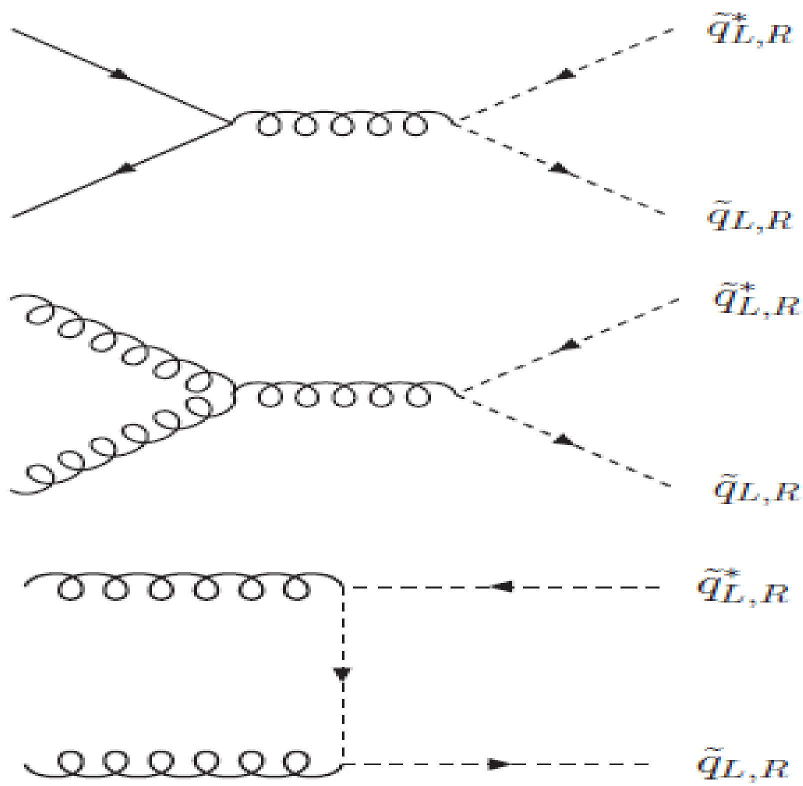
$$\delta m_{H_u}^2 = -\frac{3\lambda_t^2}{8\pi^2} m_{\tilde{t}}^2 \ln \frac{m_D^2}{m_{\tilde{t}}^2}$$

⇒ a 5 TeV Dirac gluino mass is as natural as a 900 GeV Majorana gluino in MSSM.

Dirac gluino models have several specific implications for colored sparticle production and flavor physics

Colliders:

- Gluino pair production, gluino/squark production negligible.
- t-channel Dirac gluino exchanges are mass-suppressed.
- Processes like $pp \to \tilde{q}_L\tilde{q}_L, \tilde{q}_R\tilde{q}_R$ are absent in Dirac case



(from Kribs-Martin, arxiv:1203.4821[hep-ph])

FIG. 2. The dominant tree level Feynman diagrams for squark production at the LHC in the SSSM. Dirac gluino t-channel exchange diagrams (not shown) are suppressed by $1/M_3^2$ and thus negligible. In the MSSM, by contrast, Majorana gluino exchange is suppressed by $1/\tilde{M}_3$, and thus relevant even when $\tilde{M}_3$ is large, as shown in Fig. 3.

In MSSM, the strongest FCNC constraint comes from the Kaon system (heavy gluino limit)

$$H_{\rm eff} \sim \frac{\delta_{LL}^{d,12}\delta_{RR}^{d,12}}{M^2}(\bar{d}_R s_L)(\bar{d}_L s_R)$$

In the Dirac case, there is an additional suppression

$$H_{\rm eff} \sim \delta_{LL}^{d,12}\delta_{RR}^{d,12}\frac{m_{\tilde{q}}^2}{m_D^4}(\bar{d}_R s_L)(\bar{d}_L s_R)$$

due to R-symmetry ⟹ FCNC constraints are relaxed significantly.

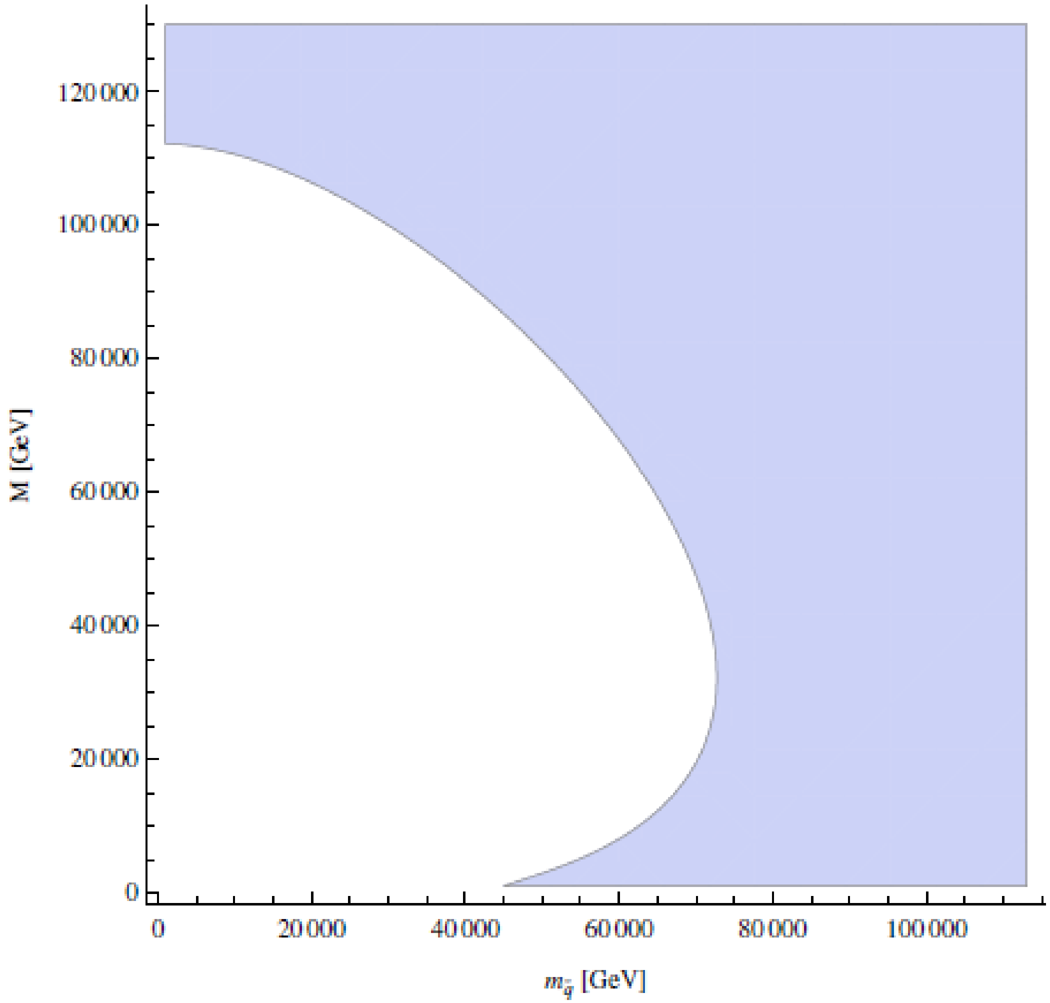


Figure 1: Contour plot in parameter space $m_{\tilde{q}}$ - M for purely Majorana gluino ($M_D = M_\chi = 0$). Along the contour $\Delta M_K = \Delta M_K^{exp} = 3.484 \times 10^{-15}$ GeV. $\sqrt{|\mathrm{Re}\,\delta^2_{12LL}|} = \sqrt{|\mathrm{Re}\,\delta^2_{12RR}|} = 0.22$.

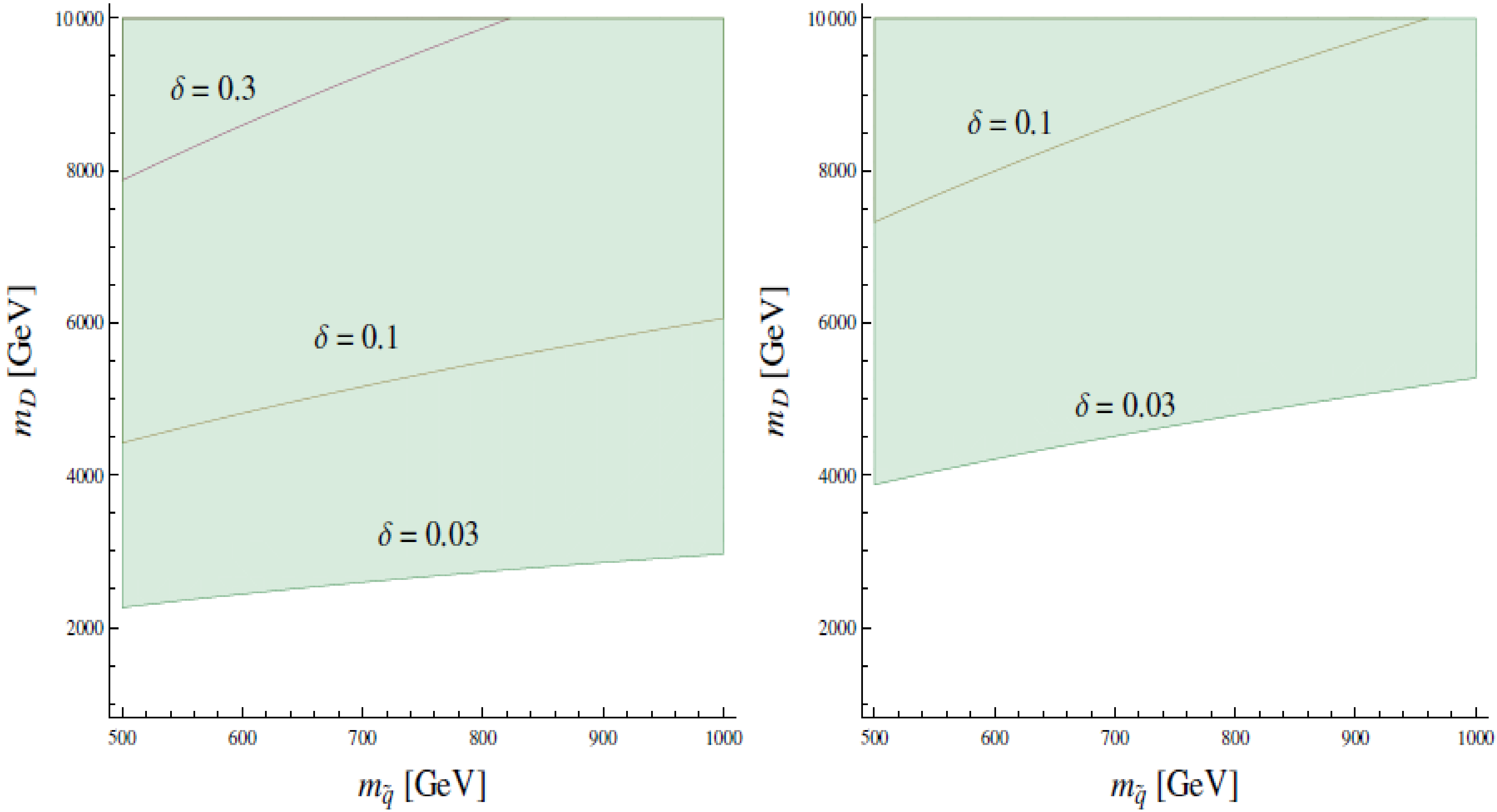


Figure 1: Contour plots in parameter space $m_{\tilde{q}}$ - m_D for purely Dirac gluino ($M = M_\chi = 0$). Left: $\delta_{LL} = \delta_{RR} = \delta$, $\delta_{LR} = \delta_{RL} = 0$. Right: $\delta_{LL} = \delta_{RR} = \delta_{LR} = \delta_{RL} = \delta$. Along the contours $\Delta M_K = \Delta M_K^{\rm exp}$ (for $\delta_{IJ} = \sqrt{|{\rm Re}\, \delta^2_{12IJ}|}$) and $\epsilon_K = \epsilon_K^{\rm exp}$ (for $\delta_{IJ} = \sqrt{c}\sqrt{|{\rm Im}\, \delta^2_{12IJ}|}$).

with c = 640

We considered a general case with:

- Majorana masses M (gluino) and M_χ (adjoint fermion)
- Dirac mass m_D
- General soft masses for scalar octet O : m_O, B_O

There are two cases with suppressed FCNC and production in colliders:

- Mostly Dirac case : $m_D >> M, M_\chi$

Protection guaranteed by R-symmetry; squarks naturally light for m_D around 5-10 TeV

- « Wrong gluino » case : $M >> m_D, M_\chi$

Lightest adjoint fermion has a small coupling to quark/squarks $g\frac{m_D}{M}$. Squarks naturally heavier than Dirac case.

One can contemplate a « wrong split SUSY scenario » with :

- Very heavy gluino, squarks and scalar octet + 1 Higgs

$$M, m_O, m_{\tilde{q}} \sim 10^{12}\ GeV$$

⇒ gauge coupling unification around $6 \times 10^{17}\ GeV$
(Bachas,Fabre,Yanagida)

- Light « wrong gluinos » + higgsinos + 1 Higgs + « wrong electroweakinos »

- The outcome is similar to split SUSY, but the light adjoint fermions are not N=1 partners (but N=2) of gauge fields !

- Lifetime of the « wrong gluinos » longer than in split SUSY.

- Different effective couplings: higgs-higgsinos-« wrong winos » vertex not anymore a gauge coupling, multiplied by m_D/M

- If by accident/tuning, a squark is in the TeV range, its low-energy effects (FCNC,production) are still strongly suppressed due to its small coupling to the light wrong gluino .

Starting with the hierarchy $m_D = \epsilon M$ and $B_{\mathcal{O}} = \epsilon^2 M^2$ radiative corrections are of the form

$$
\begin{aligned}
\delta M_\chi &\sim \epsilon^2 \frac{g_s^2}{16\pi^2} M \,, \\
\delta B_{\mathcal{O}} &\sim m_D^2 \frac{g_s^2}{16\pi^2} \log\left(\frac{\Lambda}{M}\right) \sim \epsilon^2 \frac{g_s^2}{16\pi^2} M^2 \log\left(\frac{\Lambda}{M}\right) \\
\delta m_{\mathcal{O}}^2 &\sim \frac{g_s^2}{16\pi^2} M^2 \,.
\end{aligned}
$$

A numerical example of resulting masses is

$$M \sim 10^{12}\mathrm{GeV} \gtrsim m_{\mathcal{O}} \quad \gg \quad m_D,\ \sqrt{B_{\mathcal{O}}} \sim 10^{8}\mathrm{GeV} \quad \gg \quad M_{\chi} \sim 1\mathrm{TeV}$$

➡ radiatively table to have a very heavy gluino mass compared to the « wrong gluino »

Some simple flavor models we are considering:

- One U(1) models with alignment; ex. charges

$$Q = (3,2,0) \qquad u = (3,1,0) \qquad d = (3,2,2)$$

Squark mass matrices are

$$\mathcal{M}^2_{d_L} \sim M_F^2 \begin{pmatrix} 1 & \epsilon & \epsilon^3 \\ \epsilon & 1 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}, \qquad \mathcal{M}^2_{d_R} \sim M_F^2 \begin{pmatrix} 1 & \epsilon & \epsilon \\ \epsilon & 1 & 1 \\ \epsilon & 1 & 1 \end{pmatrix}$$

and quark rotations are

$$U_L^d \sim \begin{pmatrix} 1 & \epsilon & \epsilon^3 \\ \epsilon & 1 & \epsilon^3 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix}, \qquad U_R^d \sim \begin{pmatrix} 1 & \epsilon & \epsilon \\ \epsilon & 1 & 1 \\ \epsilon & 1 & 1 \end{pmatrix}$$

There are large right-handed rotations → strongest constraints from $\tilde{Q}_1 = (\bar{d}_R \gamma_\mu s_R)^2$.

i) Pure Dirac case $m_D >> m_{\tilde{q}}$; we get

$$\frac{\langle K^0 | H_{eff} | \overline{K}^0 \rangle}{\Delta m(\exp)} \simeq 1 \times \left(\frac{\alpha_s}{0.1184} \right)^2 \left(\frac{15 \text{ TeV}}{m_{D3}} \right)^2$$

which for ϵ_K would need $m_D > 350\ TeV$, whereas for Δm_K we need $m_D > 15\ TeV$ → still very constraining.

ii) « Wrong gluino » case : for lightest octet of 5 TeV, we need Majorana gluino mass of about 250 TeV to satisfy Δm_K

In the U(1) x SU(2) model with pure Dirac gluinos, the flavor suppression is bigger than in the MSSM case, as expected.

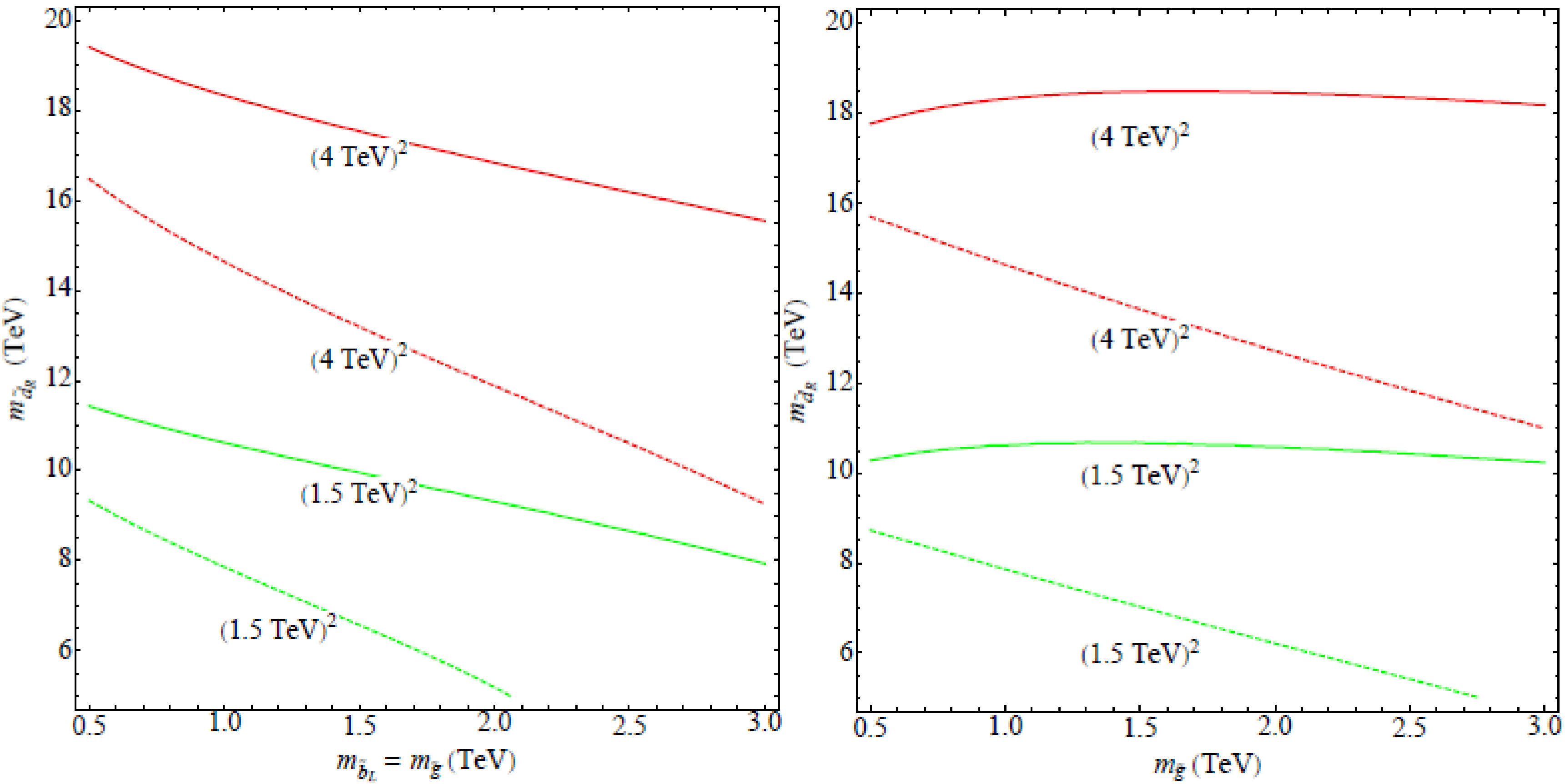


Figure 7: Constraints on the model of section 4.2.4. The dashed lines correspond to exactly Dirac gauginos, while the solid lines are purely Majorana, as in the original model of [15]. In the left plot, the left-handed sbottom mass is set equal to that of the gluino; in the right plot, the left-handed sbottom is fixed at 1 TeV.

Some string comments:

- Natural SUSY/Inverted hierarchy in string theory
- Anomalous U(1)'s in all string theories and F-theory, flavor dependent + additional discrete symmetries
- Different localization of the third generation versus the first two ones: twisted/untwisted fields, varying fluxes
- Some recent attempts to compute flavor structure of soft terms (Blumenhagen,Deser,Lust; Camara,E.D.,Palti; Camara,Ibanez,Valuenzuela).
- Dirac gauginos are natural in intersecting brane models (Antoniadis,Benakli,Delgado,Quiros and Tuckmantel)

Inverted hierarchy can also be realized in field theory:
- SUSY(SUGRA) RS 5d warped models
- flavored (higgsed) gauge mediation.

Conclusions

- Popular SUSY models are more fine-tuned, more stringent limits on SUSY spectra from direct LHC searches and flavor physics constraints.

- But there is no reason to reduce low-energy SUSY to its simplest examples: mSUGRA,CMSSM, mGMSB.

- Most theories of fermion masses generate flavor-dependent soft terms. Inverted hierarchy/natural SUSY arises naturally in Xtra dims. and string theory constructions.

- FCNC strongly constrain the flavor structure of soft terms. In MSSM, probably necessary to combine ingredients from abelian and non-abelian discrete flavor symmetries.

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- Dirac gaugino models can suppress FCNC due to

- R-symmetry in the pure Dirac gaugino mass case

- suppressed « wrong gluino » couplings to quarks/squarks in the large Majorana gluino mass case.

- Flavor models: easier to satisfy FCNC with Dirac gauginos, but still nontrivial constraints for natural values of the Dirac mass. Different collider signatures.

- Interesting to work out detailed predictions for B, D physics of flavor models.

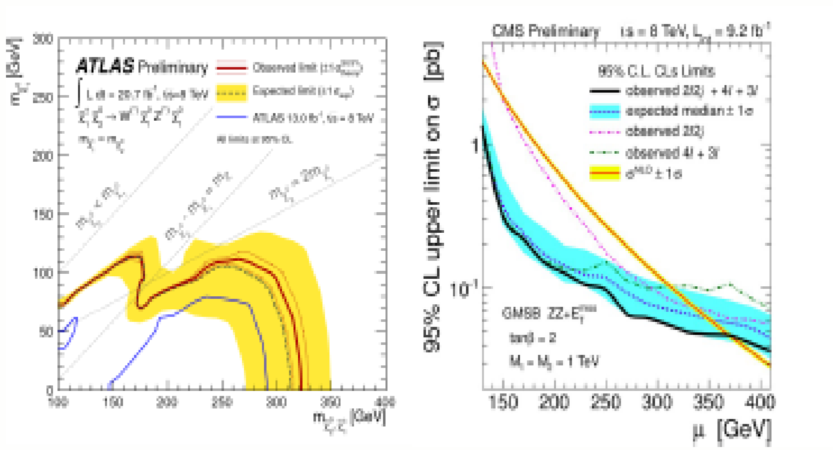
- The mechanism and the scale of SUSY is THE big unknown: split or even high-energy SUSY possible in string theory.

THANK YOU

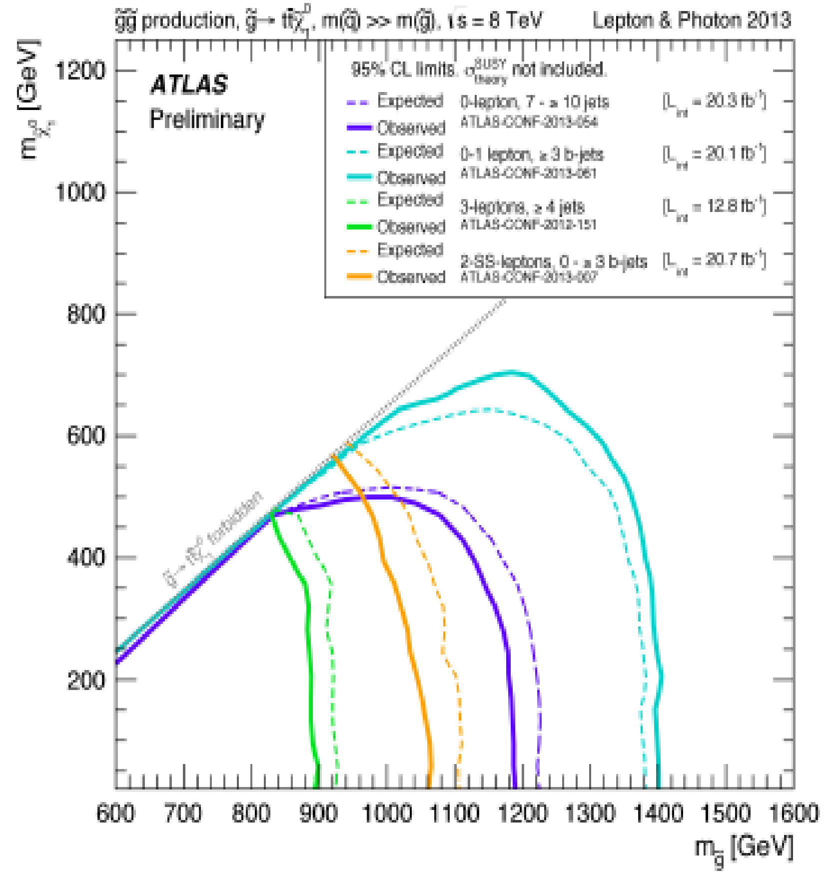
Backup slides

Bounds on « Natural SUSY » models

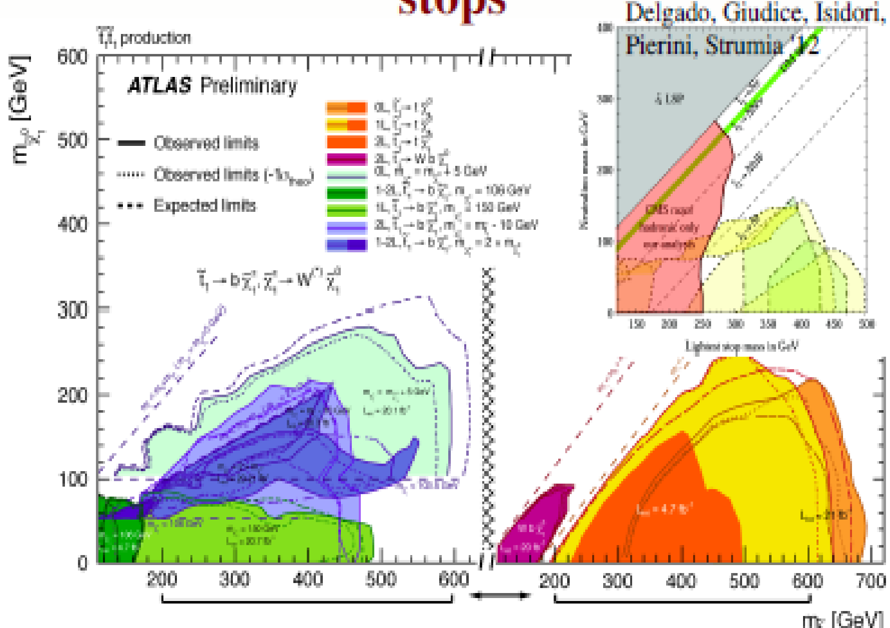
EWinos



gluino



stops



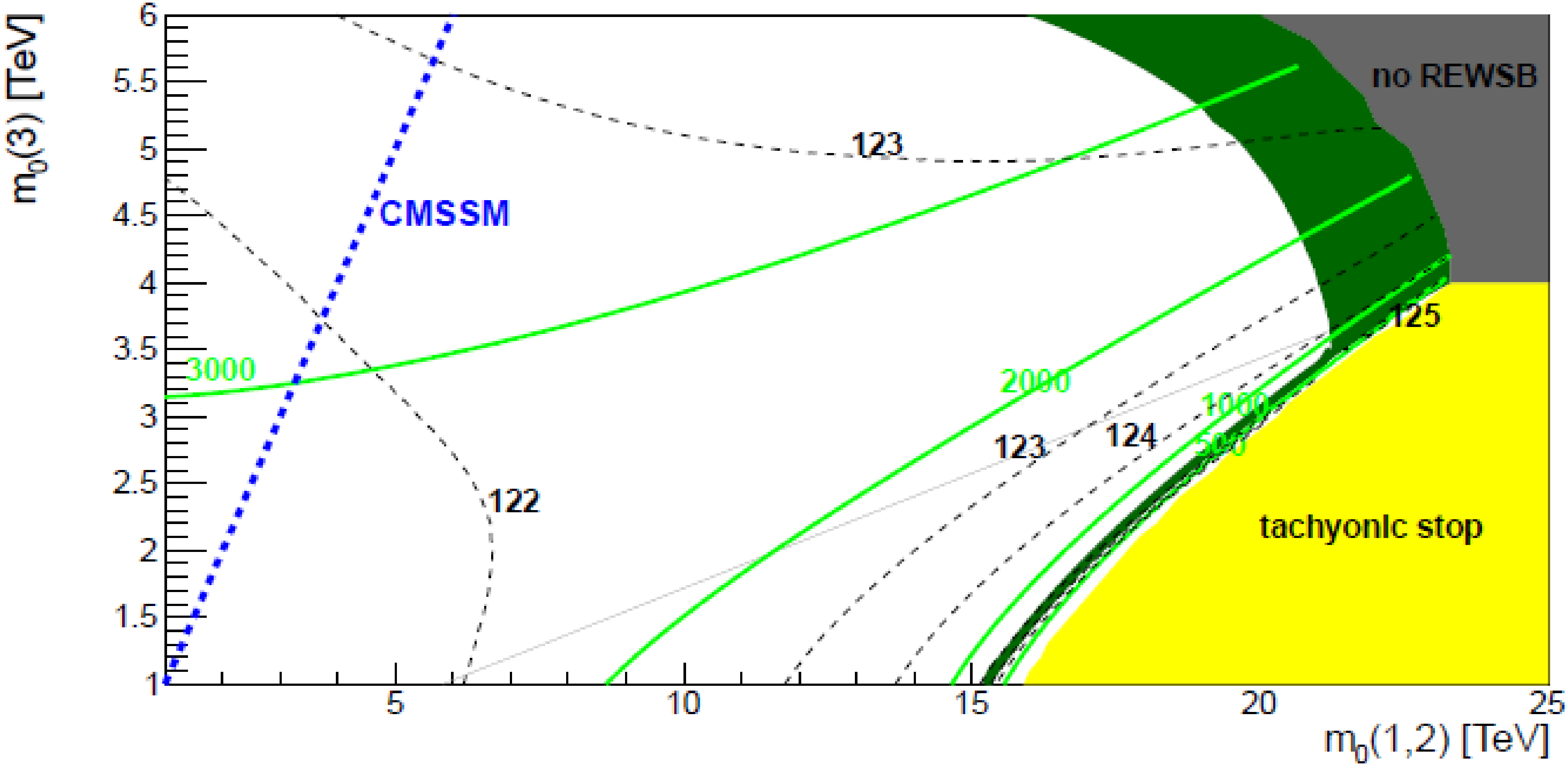


Figure 5: The same as in Figure 3 but for $M_{1/2} = 1.5$ TeV and $A_0 = 0$.

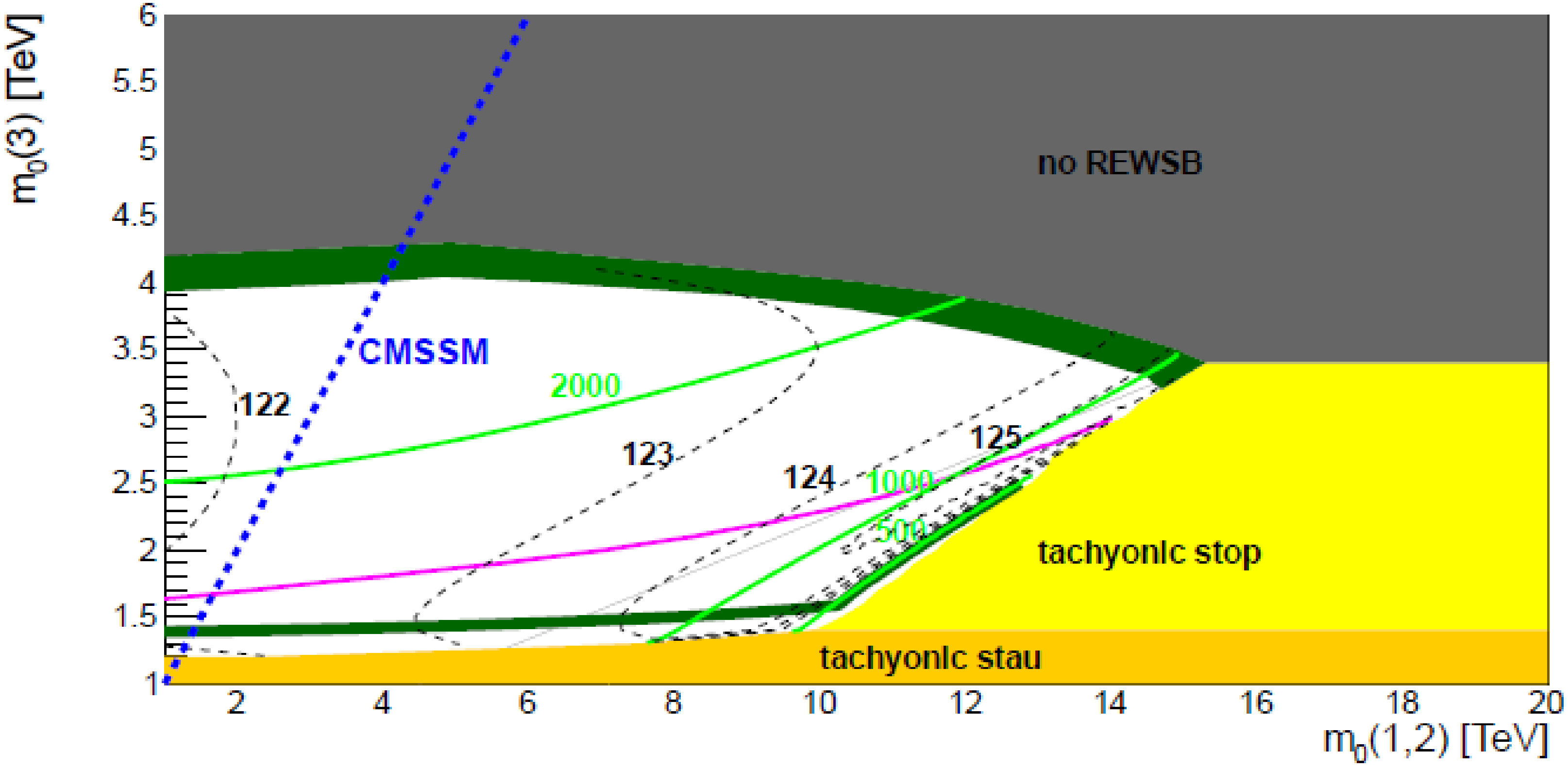
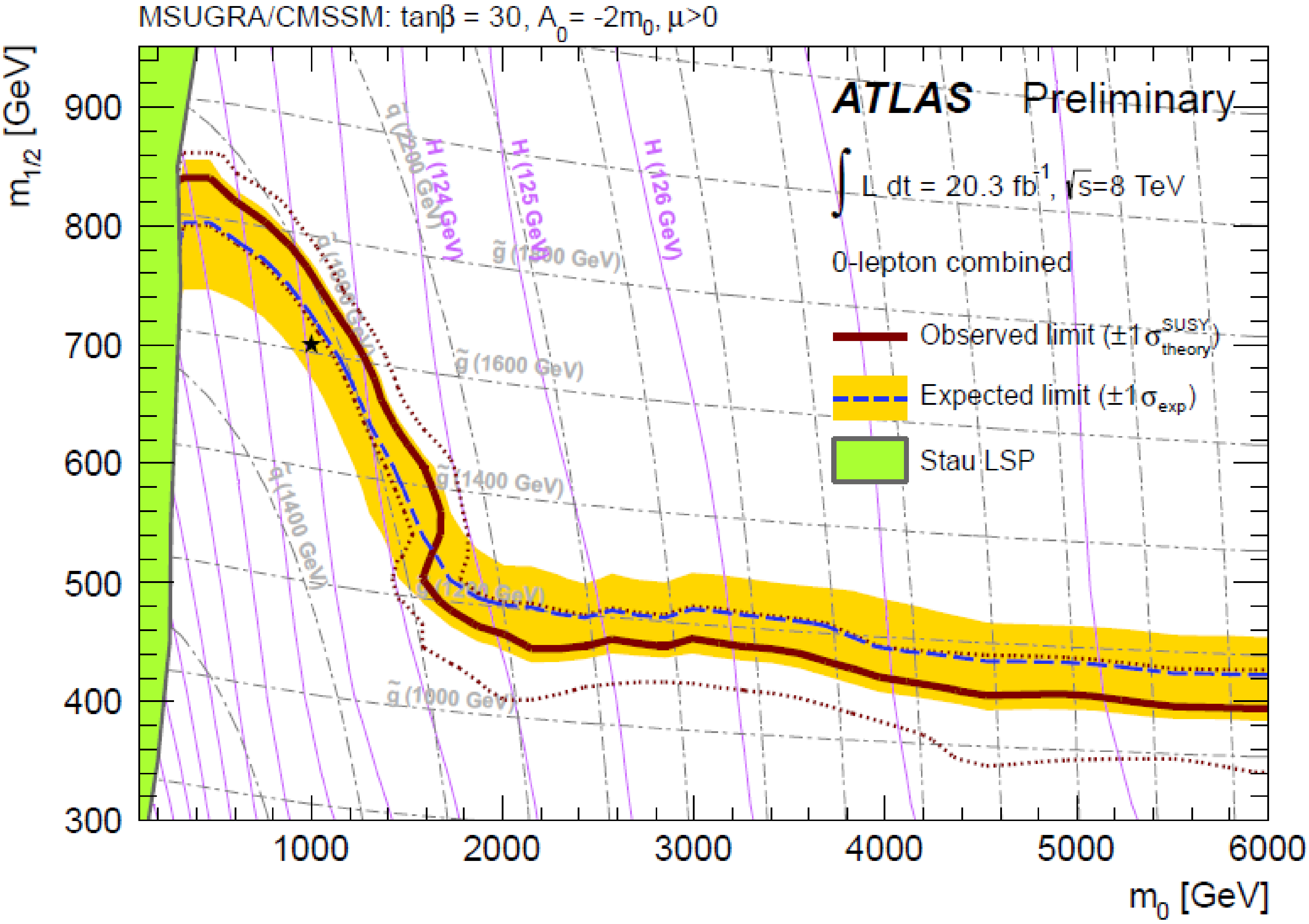


Figure 7: The same as in Figure 3 but for $\tan\beta = 50$ and $m_{H_d} = 1.6 m_0(3)$. The region below the purple line is excluded by $BR(B_s \to \mu^+\mu^-)$ at 95% C.L. The orange region is excluded because it predicts a tachyonic stau.



MSUGRA/CMSSM: tanβ = 30, A0= -2m0, μ>0
ATLAS Preliminary
∫ L dt = 20.3 fb-1, √s=8 TeV
0-lepton combined
Observed limit (±1 σ SUSY theory)
Expected limit (±1 σexp)
Stau LSP
H (124 GeV)
H (125 GeV)
H (126 GeV)
g̃ (1000 GeV)
g̃ (1200 GeV)
g̃ (1400 GeV)
g̃ (1600 GeV)
g̃ (2200 GeV)
m1/2 [GeV]
m0 [GeV]
300
400
500
600
700
800
900
1000
2000
3000
4000
5000
6000

Supersymmetry and naturalness

The **hierarchy problem** (mis?)guided BSM physics for the last 30 years.

$$\delta m_h^2 \simeq \frac{3\Lambda^2}{8\pi^2 v^2}(4m_t^2 - 4M_W^2 - 2M_Z^2 - m_h^2)$$